

Electronic Properties of the Materials
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Lecture - 12
Quantum Free Electrons: Sommerfeld Theory (Part 2)

Hello friends, in this lecture, we continue our discussion on quantum free electrons and Sommerfeld theory.

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k-space: key points to remember

- Solved TISE for free particle in a cubic box: $-\frac{\hbar^2}{2m}\nabla^2\psi(\vec{r}) = \varepsilon\psi(\vec{r})$
- Eigenfunction: $\psi_{\vec{k}}(\vec{r}) = \frac{1}{\sqrt{V}}\exp(i\vec{k}\cdot\vec{r})$, where $\vec{k}$ is wave vector
- PBC leads to discrete/quantized $\vec{k}$ values
- ▶ $k_x = \frac{2\pi n_x}{L}, k_y = \frac{2\pi n_y}{L}, k_z = \frac{2\pi n_z}{L}$
- ▶ Constructed 3D k -space with values of k_x, k_y, k_z
- ▶ Volume occupied by each point in k -space: $\Delta k = \frac{8\pi^3}{V}$
- Filling energy levels with electrons $\equiv$ filling k -space with electrons
- Pauli exclusion principle must be obeyed during filling
- Filled portion of the k -space is a sphere of radius k_F : Fermi sphere
- k_F (known as Fermi wave vector) is related to an experimentally measurable quantity: electron density $n = \frac{k_F^3}{3\pi^2}$
- Quantum free electrons at 0K have 100x kinetic energy than classical free electrons at 300K because of Pauli exclusion principle

The diagram shows a 3D grid of points in k -space, with a sphere of radius k_F centered at the origin, representing the Fermi sphere. The axes are labeled k_x, k_y, k_z .

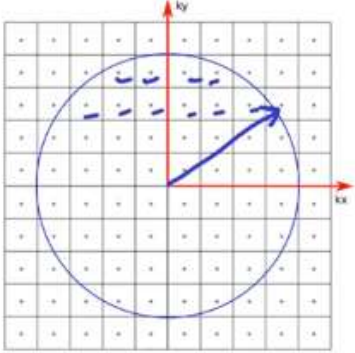
Let us briefly revise what we have done so far. So, we have solved time independent Schrodinger equation for free particle in a cubic box. And then we had bounded Eigenfunctions in this form exponential $i\vec{k} \cdot \vec{r}$ where $\vec{k}$ is the wave vector. Now, we applied periodic boundary conditions and that gave us these discrete or quantized $\vec{k}$ values where $k_x = 2\pi n_x / L$ and so on here n_x, n_y, n_z these are integers.

Now, using these discrete or quantize $\vec{k}$ values, we constructed a 3D k space with allowed values of k_x, k_y and k_z . Volume occupied by each of these discrete points is given by $\Delta k = 8\pi^3 / V$, then we fill the energy levels and then we found that filling of energy levels with electrons is same as filling of k space k electrons. During the filling Pauli exclusion principle must be obeyed that means, we cannot put more than 2 electrons per energy level or per k point. Filled portion of the k space is a sphere of radius k_F .

And this is known as the Fermi sphere k_F is known as the Fermi wave vector. And this is related to an experimentally measurable quantity that is the electron density. Electron density $n = k_F^3 / 3\pi^2$. We also found that one term free electrons at 0 Kelvin have 100 times more kinetic energy than classical free electrons at 300 Kelvin and this happens because of the Pauli exclusion principle.

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Thermodynamic properties at $T = 0$ K



- Let $F(\vec{k})$ be a smooth function of $\vec{k}$
- Example: $\varepsilon(\vec{k}) = \frac{\hbar^2 k^2}{2m}$ is a smooth quadratic function of $\vec{k}$
- To get total energy we have to calculate the sum: $\sum_{k < k_F} \varepsilon(\vec{k})$
- The sum $\sum_k$ can be converted to integral

$\sum_k F(\vec{k}) = \frac{V}{8\pi^3} \sum_k F(\vec{k}) \Delta k$, where $\Delta k = 8\pi^3/V$ is volume of

Let us calculate some thermodynamic properties at ground state. Say we know some property as a function of k and we have a smooth function of k . For example, energy = $\hbar^2 k^2 / 2m$ is a smooth function of k . Now, we want to calculate the total energy of the system. Now, this is the Fermi sphere with radius k_F and the value of k_F is given to us. To calculate the total energy, we have to add up the energies of all the k states lying inside the Fermi sphere. For example, we know that all these states, these are lying inside the Fermi sphere.

So, we have to add the energy of all the states inside the Fermi sphere and then we can get the total energy. So, we have to calculate this sum. Let us convince ourselves that the sum can be converted to an integral under certain condition.

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• To get total energy we have to calculate the sum: $\sum_{k < k_F} \epsilon(\vec{k})$

• The sum $\sum_k$ can be converted to integral

• $\sum_k F(\vec{k}) = \frac{V}{8\pi^3} \sum_k F(\vec{k}) \Delta k$, where $\Delta k = 8\pi^3/V$ is volume of k -space per allowed $\vec{k}$ value

• If $V \rightarrow \infty / \Delta k \rightarrow 0$ and $F(\vec{k})$ varies slowly, sum converted to integral

• $\lim_{V \rightarrow \infty} \frac{1}{V} \sum_{k < k_F} F(\vec{k}) = \int \frac{d\vec{k}}{8\pi^3} F(\vec{k})$

$\frac{1}{\Delta k} \sum_k F(\vec{k}) \Delta k = \frac{V}{8\pi^3} \sum_k F(\vec{k}) \Delta k$

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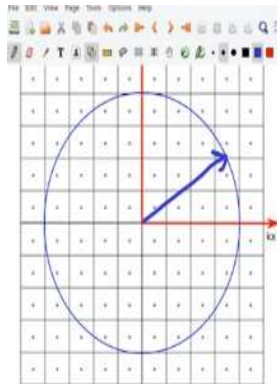
So, we have to calculate this sum for all the allowed values of k . Now, we can take this so sum over k F of k we want to calculate the sum. So

$$\frac{1}{\Delta k} \sum_k F(\vec{k}) \Delta k = \frac{V}{8\pi^3} \sum_k F(\vec{k}) \Delta k$$

Now, we use some limit V goes to infinity, then you see that this box they become Δk becomes very small.

So, each of these boxes it will become very small as V goes to infinity and then this discrete K points they will become almost continuous and that is why in this limit we get convert this sum to an integral so, what we do that in this equation correct being V to the left hand side so, that means, this left hand side will become 1 by V sum over k less than k_F of k . and then the right hand side we have 1 by $8\pi^3$ cube. So, we have this term and we have Δk is now, we just write it as dk and then we have this function F of k and then as I said that this discrete space becomes almost continuous in the limit ΔV going to infinity or Δk goes to 0. So, that means, we just we can do this calculate this integral.

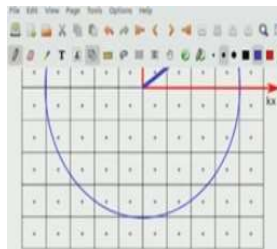
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- To get thermodynamic property, we have to calculate the sum: $\sum_k F(\vec{k})$
- The sum $\sum_k$ can be converted to integral
- $$\lim_{V \rightarrow \infty} \frac{1}{V} \sum_{k < k_F} F(\vec{k}) = \int \frac{d\vec{k}}{8\pi^3} F(\vec{k})$$
- Let us derive total # electrons using the above recipe
- In this case, $F(\vec{k}) = 1$ and $d\vec{k} = 4\pi k^2 dk$
- $$\frac{N}{V} = 2 \int_0^{k_F} \frac{4\pi k^2 dk}{8\pi^3} = \frac{k_F^3}{3\pi^2}$$

Let us derive total number of electrons using the above method; we have to find out how many states are there within the Fermi sphere. And then we have to multiply with 2 as there are 2 electrons per state according to the Pauli's exclusion principle that will give us the total number of electrons in the system. In this case, the function the smooth function is F of k this is equal to 1 and then we write the volume element in the case space dk as 4 pi k square dk we can write it in this form, because we have spherical symmetry. Now, let us calculate this.

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- $$\lim_{V \rightarrow \infty} \frac{1}{V} \sum_{k < k_F} F(\vec{k}) = \int \frac{d\vec{k}}{8\pi^3} F(\vec{k})$$
- Let us derive total # electrons using the above recipe
- In this case, $F(\vec{k}) = 1$ and $d\vec{k} = 4\pi k^2 dk$
- $$\frac{N}{V} = 2 \int_0^{k_F} \frac{4\pi k^2 dk}{8\pi^3} = \frac{k_F^3}{3\pi^2}$$
- $$N = 2 \sum_{k=0}^{k_F} F(\vec{k})$$
- $$\frac{N}{V} = \frac{2}{V} \sum_{k=0}^{k_F} F(\vec{k}) = 2 \int \frac{d\vec{k}}{8\pi^3} \cdot 1$$

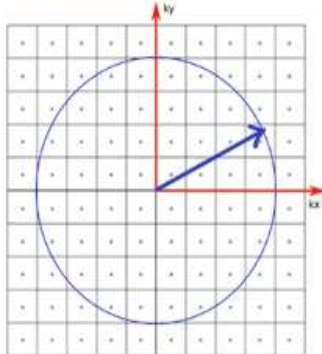
So, total number of electrons $N = 2 \sum_{k=0}^{k_F} F(\vec{k})$. Now, we multiply to because of Pauli's exclusion principle. So, this sum will give you the total number of steps within the Fermi sphere and then

we have to multiply with 2 because each state can have 2 electrons and now, we divided by volume such that $\frac{N}{V} = \frac{2}{V} \sum_{k=0}^{k_F} F(\vec{k})$ and I know that 1 by V sum over k less then F of k this is given by this integral. So, that means, N by V can be better than in this form.

So, I just replaced this sum with the $2 \int \frac{d\vec{k}}{8\pi^3}$ and then I know that this f of k=1 and then in place of begin I just replaced 4 pi k square dk. So, this is the integral that we have to evaluate and this is not so, difficult to evaluate. So, if we just do this, so, this is like 1 over pi square integral 0 to k F k square dk and in that case, we see that N by V = k cube F divided by 3 pi square.

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Thermodynamic properties at $T = 0$ K



- Let us derive total energy E of electron gas
- $E = 2 \sum_{k < k_F} \frac{\hbar^2 k^2}{2m}$, thus $F(\vec{k}) = \hbar^2 k^2 / 2m$
- The sum $\sum_k$ can be converted to integral
- $$\lim_{V \rightarrow \infty} \frac{1}{V} \sum_{k < k_F} F(\vec{k}) = \int \frac{d\vec{k}}{8\pi^3} F(\vec{k})$$
- $$\frac{E}{V} = 2 \int_0^{k_F} \frac{4\pi k^2 dk}{8\pi^3} \frac{\hbar^2 k^2}{2m} = \frac{\hbar^2 k_F^5}{10\pi^2 m}$$

Let us derive total energy of the electron gas again the Fermi sphere is given to us the radius of the sphere. So, what we have to do is that, we have to find all the states lying within this Fermi sphere and then we have to add the energy of every state to get the total energy. So, that means we have to calculate this sum. And then we have F of k the function that we have to integrate is h cross square k square divided by 2m known that this factor to these accounts for the Pauli exclusion principle that means we have 2 electrons, but state now, let us calculate this.

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$$\lim_{V \rightarrow \infty} \frac{1}{V} \sum_{k < k_F} F(\vec{k}) = \int \frac{dn}{8\pi^3} F(\vec{k})$$

$$\frac{E}{V} = \frac{2}{V} \int_0^{k_F} \frac{4\pi k^2 dk}{8\pi^3} \frac{\hbar^2 k^2}{2m} = \frac{\hbar^2 k_F^5}{10\pi^2 m}$$

$$\frac{N}{V} = \frac{k_F^3}{3\pi^2}$$

• Energy per particle: $\frac{E}{N} = \frac{3}{5} \frac{\hbar^2 k_F^2}{2m} = \frac{3}{5} \epsilon_F = \frac{3}{5} k_B T_F \sim 10^4 K$

• Compare with classical electron gas: $\frac{E}{N} = \frac{3}{2} k_B T \rightarrow \sim 300 K$

• Remember that $T_F \sim 10^4 K$

$$\frac{E/N}{N/V} = \frac{3}{5} \frac{\hbar^2 k_F^2}{2m} = \frac{3}{5} \epsilon_F$$

$$\frac{E}{V} = \frac{1}{\pi^2} \frac{\hbar^2}{2m} \int_0^{k_F} k^4 dk = \frac{\hbar^2 k_F^5}{10\pi^2 m}$$

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So, we know that the sum can be converted to some integral of this form and then we also know that this is equal to E. So, this is where we are just adding the energies of all the states. So, this will give E. So, that means, this factor on the left-hand side that is equal to E by V and then we have this factor of 2 which accounts for that Pauli exclusion principle and in the right-hand side we have this integral and then we know that $dk = 4\pi k^2 dk$ and this is because of the spherical symmetry and then we have this function f of $k = \hbar^2 k^2 / 2m$.

So, you have to calculate this integral and then this integral will be equal to

$$\frac{1}{\pi^2} \frac{\hbar^2}{2m} \int_0^{k_F} k^4 dk = \frac{\hbar^2 k_F^5}{10\pi^2 m}$$

So, now, what we know is we know that $N/V = k_F^3 / 3\pi^2$ and we know that this term is equals to E/V . E/V is equals to this term.

So, that means, we can also write this ratio of E/V and then N/V and that will turn out to be $3/5$ th of $\hbar^2 k_F^2 / 2m$ and then we know that this $\hbar^2 k_F^2 / 2m = E_F$ And the Fermi of the square $3/5$ of E_F and then E_F we can convert it to some temperature scale we can just write E_F is equals to Boltzmann constant k_B times T_F where T_F is some temperature scale.

$$\frac{E/V}{N/V} = \frac{3}{5} \frac{\hbar^2 k_F^2}{2m} = \frac{3}{5} \epsilon_F$$

Now, we can compare this with classical the energy of the classical electron gas and that turns out to be 3 by 2 times K B T whereas, this is of the order of 10 power for Kelvin and the room temperature T this is of the order of 300 Kelvin. Thus, we again see that the quantum free electron gas has almost 100 times more energy than the classical free electron gas.

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Thermodynamic properties at $T = 0$ K

- Combined 1st and 2nd law of thermo: $dE = TdS - PdV + \mu dN$
- Pressure exerted by free electron gas at $T = 0$ K: $P = - \left(\frac{\partial E}{\partial V} \right)_{N,S}$

$$E = \frac{3}{5} N \epsilon_F = \frac{3}{5} N \frac{\hbar^2 k_F^2}{2m} \quad \left| \quad k_F = \left(\frac{3\pi^2 N}{V} \right)^{1/3} \right.$$

$$E = \frac{3N\hbar^2}{10m} \frac{(3\pi^2 N)^{2/3}}{V^{2/3}} = C \times V^{-2/3}$$

- E increases as V decreases: $E \propto V^{-2/3}$

We already have calculated the internal energy of free electron gas, let us calculate few more thermodynamic properties combining first and second law of thermodynamics we can write dE so, this is the E is the internal energy equal to TdS, T is the temperature, S is the entropy minus PdV the P is pressure, V is volume plus mu dN mean is chemical potential and N is the total number of particles. Now, if we take this equation, then we can calculate the pressure exerted by the free electron gas that T = 0 Kelvin.

So, by calculating this partial derivative del E del V that constant N and S that is equal to pressure.

Let us calculate this term. So, we already know that $E = \frac{3}{5} N \epsilon_F$ where N is the total number of particles, total number of electrons and E F is the Fermi energy. Now, we know that the Fermi energy can further written as $\frac{3}{5} N \frac{\hbar^2 k_F^2}{2m}$. Now, we also know that

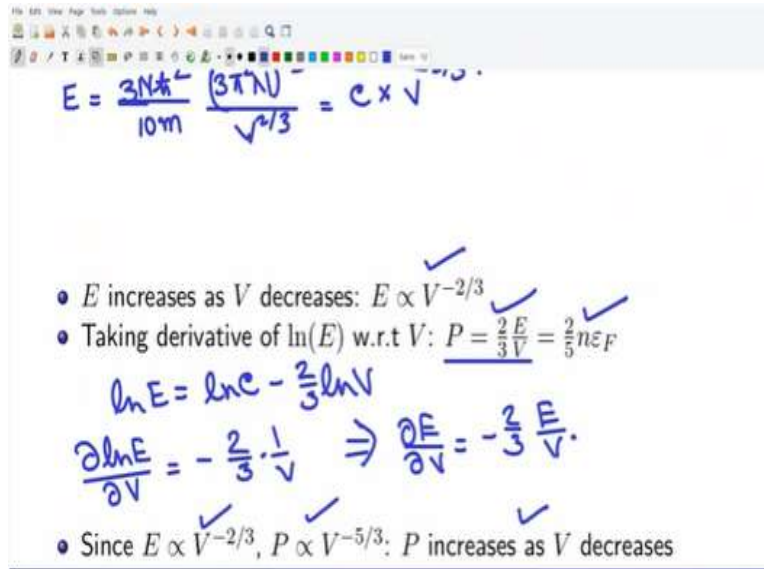
$$k_F = \left(\frac{3\pi^2 N}{V} \right)^{1/3}$$

So, that means,

$$E = \frac{3N\hbar^2}{10m} \frac{(3\pi N)^{2/3}}{V^{2/3}} = cV^{-2/3}$$

so, thus energy increases as V decreases. So, energy is proportional to V power -2 by 3.

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The image shows a handwritten derivation of the Fermi energy equation and its derivative. The equation is written as $E = \frac{3N\hbar^2}{10m} \frac{(3\pi N)^{2/3}}{V^{2/3}} = c \times V^{-2/3}$. Below this, a list of points is provided:

- E increases as V decreases: $E \propto V^{-2/3}$
- Taking derivative of $\ln(E)$ w.r.t V: $P = \frac{2}{3} \frac{E}{V} = \frac{2}{5} n \epsilon_F$

 The derivation of the derivative is shown as $\ln E = \ln c - \frac{2}{3} \ln V$, followed by $\frac{\partial \ln E}{\partial V} = -\frac{2}{3} \cdot \frac{1}{V} \Rightarrow \frac{\partial E}{\partial V} = -\frac{2}{3} \frac{E}{V}$. Finally, it states:

- Since $E \propto V^{-2/3}$, $P \propto V^{-5/3}$. P increases as V decreases

Now, we can take log of this equation and then we can write log of E = log of C - 2 third log of V and then we calculate this derivative, del we calculate the partial derivative del log E divided by del V = -2 by 3 into 1 by V. So, this implies that we can further write del E by del V = -2 by 3 E by V. So, this is what we have and then we can further write this term like this. Since, energy goes as V power minus 2 by 3 then from this equation write from this equation.

We can see that all we can find out that pressure goes as V power -5 by 3. So, pressure increases as volume decreases thus, as one tries to compress a metal free electron pressure is going to resist.

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• Since $E \propto V^{-2/3}$, $P \propto V^{-5/3}$: P increases as V decreases

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Thermodynamic properties at $T = 0$ K

- Bulk modulus (inverse of compressibility): $B = -V \left(\frac{\partial P}{\partial V} \right)$
- Taking derivative of $\ln(P)$ w.r.t. V : $B = \frac{5}{3}P = \frac{10}{9}E = \frac{2}{3}n\varepsilon_F$

$$\ln P = \ln C - \frac{5}{3} \ln V$$

$$\frac{\partial \ln P}{\partial V} = -\frac{5}{3} \cdot \frac{1}{V} \Rightarrow \frac{\partial P}{\partial V} = -\frac{5}{3} \frac{P}{V}$$

n (cm^{-3})	k_F (cm^{-1})	v_F (cm s^{-1})	ε_F (eV)	T_F (K)

The resistance to compression is known as the bulk modulus. It is defined as the ratio of infinitesimal pressure increase due to the decrease of volume. Let us calculate these all. So, again we get log such that we can write log of P is equals to, so, we are just taking log of this equation so, $\log$ of $P = \log$ of some constant minus 5 by 3 $\log$ of V and again we calculate the derivative points which is equal to -5 by 3 into 1 by V and then we can write this in this form $\frac{\partial P}{\partial V} = -5 \text{ by } 3 \frac{P}{V}$.

Thus we can write bulk modulus B is equals to $-V \cdot \frac{\partial P}{\partial V}$. So, minus V will cancel this V and then $\frac{\partial P}{\partial V}$. So, it is like 5 by 3 of P . So, that means bulk modulus will be 5 by 3 times P and then we can further express it in terms of energy as well as Fermi energy.

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• Taking derivative of $\ln(P)$ w.r.t. V : $B = \frac{5}{3}P = \frac{10}{9} \frac{E_F}{V} = \frac{2}{3} n \epsilon_F$

$\ln P = \ln C - \frac{5}{3} \ln V$
 $\frac{\partial \ln P}{\partial V} = -\frac{5}{3} \cdot \frac{1}{V} \Rightarrow \frac{\partial P}{\partial V} = -\frac{5}{3} \frac{P}{V}$

	n (cm^{-3})	k_F (cm^{-1})	v_F (cm s^{-1})	ϵ_F (eV)	T_F (K)
Li	4.70×10^{22}	1.11×10^8	1.29×10^8	4.72	5.48×10^4
Na	2.65	0.92	1.07	3.23	3.75
Cu	8.45	1.36	1.57	7.00	8.12
Ag	5.85	1.20	1.38	5.48	6.36
Be	24.2	1.93	2.23	14.14	16.41
Mg	8.60	1.37	1.58	7.13	8.27
Al	18.06	1.75	2.02	11.63	13.49

So, that means, you see the bulk modulus can be expressed in terms of the electron density as well as the Fermi energy, but now, note that we already have estimated the value of n electron density and the value of the Fermi energy thus, we can compare the calculated value with the experimental value.

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Thermodynamic properties at $T = 0$ K

• Let us compare with experimental values n, ϵ_F

Element	Free electron theory	Experiment
Li	23.9 GPa	11.5 GPa
Na	9.23	6.42
K	3.19	2.81
Rb	2.28	1.92
Cs	1.54	1.43
Cu	63.8	134.3
Ag	34.5	99.9
Al	228	76

In this column, we have the experimental values of bulk modulus for different metals and then in this column we have the values calculated from free electron theory. So, what we did is that we calculated the value of an electron density we calculated the value of E_F and from that, we can find the value of bulk modulus.

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K	3.19	2.81
Rb	2.28	1.92 ✓
Cs	1.54	1.43 ✓
Cu	63.8	134.3 ✓
Ag	34.5	99.9 ✓
Al	228	76

- For certain alkali metals, very good agreement ✓
- As we try to compress metal, pressure of free electron gas resists ✓
- But this can not be the only source of resistance to compression – thus mismatch expected ✓
- However, free electron gas pressure is as important as other effects ✓

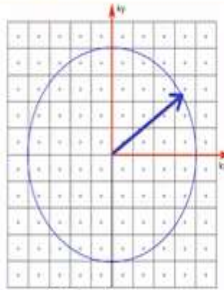
And then you will see that the values for certain alkali metals like rubidium and cesium the agreement is really good. It is not so, good for say like copper or silver and so on. So, what we can conclude from this as we try to compress a metal the pressure of free electron gas will resist, but this cannot be the only source of resistance to compression. So, the mismatch that we find here is expected.

But you see that the filling on gas pressure is as important as other effects because, if you at least look at the order of magnitude, these numbers they are of the same order of magnitude. So, that leads to the conclusion that that free electronic gas pressure is as important as other effects.

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Thermodynamic properties at $T = 0$ K: key points

- Thermodynamic variables (like N, E) calculated from sum: $\sum_{\vec{k}} F(\vec{k})$ ✓
- $F(\vec{k})$ depends on what we want to calculate ✓
 - ▶ Example: $F(\vec{k}) = 1$ if we want to calculate N ✓
 - ▶ Example: $F(\vec{k}) = \frac{\hbar^2 k^2}{2m}$ if we want E ✓
- Not easy to sum for $\sim 10^{23}$ terms ✓
- Analytical calculations simpler if we convert $\sum_{\vec{k}} \Rightarrow \frac{V}{8\pi^3} \int d\vec{k}$ ✓
- $\sum_{\vec{k}} F(\vec{k}) = \frac{V}{8\pi^3} \int d\vec{k} F(\vec{k})$ ✓



Before we go further, let us review what we have done so far. So, we have found that thermodynamic variables like the total number of electrons or the energy of the electrons can be calculated from this sum and this sum needs to be calculated up to the Fermi sphere. Now, F of k depends on what do you want to calculate for example, F of $k=1$ if we want to calculate the total number of electrons, F of k is $\hbar^2 k^2$ by $2m$ if we want to calculate the total energy and so, on.

Now, note that this is not an easy solve to calculate, because the number of electrons that we have in the system is huge it is of the order of 10^{23} . Analytical calculation it becomes simpler if we convert the sum to the integral

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▶ Example: $F(k) = 1$ if we want to calculate N
 ▶ Example: $F(k) = \frac{\hbar^2 k^2}{2m}$ if we want E

- Not easy to sum for $\sim 10^{23}$ terms
- Analytical calculations simpler if we convert $\sum_k \Rightarrow \frac{V}{8\pi^3} \int d\vec{k}$
- $\sum_k F(\vec{k}) = \frac{V}{8\pi^3} \int d\vec{k} F(\vec{k})$
 - ▶ Correct in the limit $d\vec{k} \rightarrow 0/V \rightarrow \infty$
 - ▶ In this limit, discrete points in k -space becomes *almost* continuous
- $\lim_{V \rightarrow \infty} \frac{1}{V} \sum_k F(\vec{k}) = \int \frac{d\vec{k}}{8\pi^3} F(\vec{k}) \Leftarrow$

$\Delta k = \frac{8\pi^3}{V}$

And then this is how the sum can be written as an integral and this conversion is correct in the limit Δk goes to 0 or that means in the limit V is going to infinity. So, that means, so, this is like Δk right the point and occupied other volume occupied by one point. So, we see that as V goes to ∞ , $\Delta k = 8\pi^3/V \rightarrow 0$. So, as V goes to infinity, the volume occupied by each point this will go to almost 0.

So, that means the discrete space now can be looked at as some almost continuous space and then we can convert the sum to this integral and then we can write it in this form. And we have used this to evaluate the total number of electrons as well as the total energy of the electrons.

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Thermodynamic properties at $T = 0$ K: summary

- Fermi wave vector related to electron density: $k_F = (3\pi^2 n)^{1/3}$
- Define Fermi energy: $\varepsilon_F = \frac{\hbar^2 k_F^2}{2m} = \frac{\hbar^2}{2m} (3\pi^2 n)^{2/3}$
- Ground state thermodynamic properties expressed in terms of n

Property	General expression	In terms of n, ε_F
Energy per unit volume	$\frac{E}{V}$	$\frac{3}{5} n \varepsilon_F$
Energy per electron	$\frac{E}{N}$	$\frac{3}{5} \varepsilon_F$
Pressure	$P = \frac{2}{3} \frac{E}{V}$	$\frac{2}{5} n \varepsilon_F$
Bulk modulus	$B = \frac{5}{3} P$	$\frac{2}{3} n \varepsilon_F$

- Values of n and ε_F for selected metals

	n (cm ⁻³)	k_F (cm ⁻¹)	v_F (cm s ⁻¹)	ε_F (eV)	T_F (K)
Li	4.70×10^{22}	1.11×10^8	1.29×10^8	4.72	5.48×10^4
Na	2.65	0.92	1.07	3.23	3.75

Let me summarize some of the thermodynamic properties of free electron gas at ground state. The beauty of free electron gas model is its simplicity, we need only one parameter in the electron density and from that we can calculate k_F the Fermi wave vector and from Fermi wave vector we can calculate ε_F the Fermi energy and once we have this n and ε_F , we can express all the thermodynamic properties.

For example, energy per unit volume, this is given by $\frac{3}{5} n \varepsilon_F$, energy per electron this is given by $\frac{3}{5} \varepsilon_F$, pressure is given by $\frac{2}{5} n \varepsilon_F$, bulk modulus is given by $\frac{2}{3} n \varepsilon_F$. Now, although we are using both n and ε_F , actually we have only one independent variable and because ε_F is related to n by this equation.

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• Ground state thermodynamic properties expressed in terms of n and ε_F

Property	General expression	In terms of n and ε_F
Energy per unit volume	$\frac{E}{V}$	$\frac{3}{5}n\varepsilon_F$
Energy per electron	$\frac{E}{N}$	$\frac{3}{5}\varepsilon_F$
Pressure	$P = \frac{2}{3}\frac{E}{V}$	$\frac{2}{5}n\varepsilon_F$
Bulk modulus	$B = \frac{5}{3}P$	$\frac{2}{3}n\varepsilon_F$

• Values of n and ε_F for selected metals

	n (cm^{-3})	k_F ($\text{\AA}^{-1}$)	v_F (cm s^{-1})	ε_F (eV)	T_F (K)
Li	4.70×10^{22}	1.11×10^8	1.29×10^8	4.72	5.48×10^4
Na	2.65	0.92	1.07	3.23	3.75
Cu	8.45	1.36	1.57	7.00	8.12
Ag	5.85	1.20	1.38	5.48	6.36

• Theoretically predicted and experimentally measured B : order of magnitude matches

So, now, if you look at this table, once we have found the value of n , we can find the value of ε_F and from that we can calculate all the ground state properties listed in this table. Among all the properties, we can actually compare with experimental result in case of bulk modulus. And then, we already have seen that the theoretical predicted and experimentally measured value of bulk modulus for most of the metals the order of magnitude matches. This provides us some kind of validation of quantum free electron gas model.