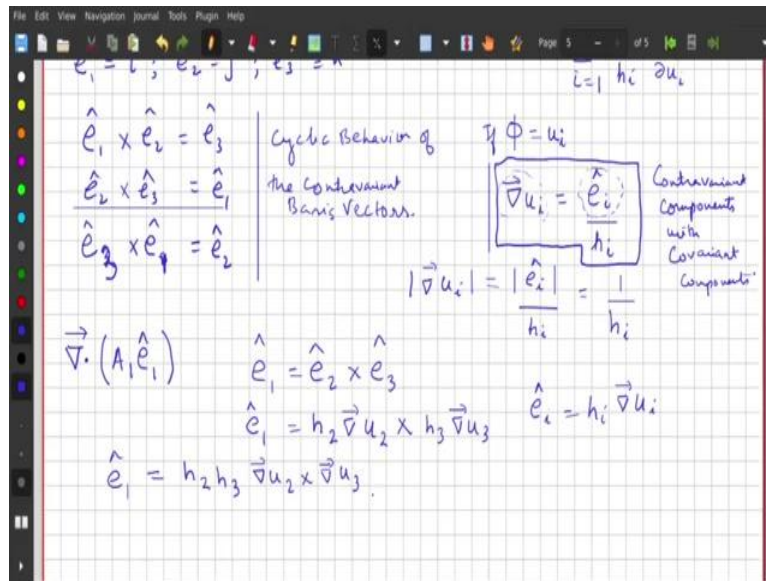


Statistical Thermodynamics for Engineers
Professor. Saptarshi Basu
Indian Institute of Science, Bengaluru
Lecture 23
Supplementary Video 8 Coordinate System 3

Welcome everyone to another segment of introductory lectures where we were discussing about coordinate systems in general.

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And in the last lecture, we saw the cyclic behavior. So, we saw the cyclic behavior of the basis vectors, cyclic behavior of the contravariant and basis vectors. So, we will continue from here in this lecture and develop the general ideas of writing the various vector operators like the grad, the divergence and the Laplacian in a general fashion in any coordinate system. So, we are developing the idea from a generic co-ordinate system, perspective that we started with basis vectors in an arbitrary curvilinear coordinate in three dimensions.

So, let us see. So, remember where we had this expression where we had grad of, let us say, this could be represented as a sum over i equals 1 to 3 e i at by the scale factor h i, partial phi, partial U i. So, let us use this important identity that we developed and using this let us try to find out what that if phi is the ith curvilinear coordinate, let us say, remember this ith curvilinear coordinate? Like, we had u1, u2 u3 in three space and we were thinking about u1 equals to c1, the surface in that three-dimensional space.

So, let us say if phi equals u i, what will this equation become? So, this equation will become grad phi, instead of phi, we have u i and this will become e i at i h i, where if we substitute phi

equal to u_i , the sum will only survive for the component where like, ϕ is u_i , so the other the cross terms will not be there because like the derivative of let us say u_1 with respect to u_2 will be 0 because u_1 and u_2 are orthogonal, and they are independent coordinates.

So, grad of u_i will become \hat{e}_i hat by h_i , where this h_i is a scale factor. So, the magnitude of grad of u_i is the magnitude of \hat{e}_i hat by h_i . So, then, this becomes, since this is a unit vectors, so this becomes 1 over h_i . So, this is a very important expression that we got and you see this relates somehow the contravariant components with the covariant components, because if you see the right hand side e serve i , this represents the contravariant components but the grad u_i , this is the normal component, so this represents the covariant basis.

So, let us keep in mind this very, this identity will become very handy as we develop further ideas. So, let us now what do we have? Let us say we want to write the divergence operator. So, let us say we are in divergence of $A_1 e_1$ vector. So, if you see, this we can write using, so let us say we use this middle expression.

If we use this middle expression, we can write e_1 as e_2 cross product e_3 and then using this expression, which we just now derived, where we had \hat{e}_i hat is h_i grad u_i , this we can rewrite as h_2 grad u_2 cross h_3 grad u_3 . So, this becomes, so this is e_1 cap. So, therefore the e_1 vector becomes $h_2 h_3$, so these are the scale factors in h_2 and h_3 multiplied by grad of u_2 cross product grad of u_3 . So, that is what e_1 is.

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$$\hat{e}_1 = \frac{h_2 \vec{\nabla} u_2 \times h_3 \vec{\nabla} u_3}{h_2 h_3 \vec{\nabla} u_2 \times \vec{\nabla} u_3}$$

$$\hat{e}_i = h_i \vec{\nabla} u_i$$

$$\hat{e}_2 = h_3 h_1 \vec{\nabla} u_3 \times \vec{\nabla} u_1$$

$$\hat{e}_3 = h_1 h_2 \vec{\nabla} u_1 \times \vec{\nabla} u_2$$

$$\vec{\nabla} \cdot (A_1 \hat{e}_1) = \vec{\nabla} \cdot (A_1 h_2 h_3 \vec{\nabla} u_2 \times \vec{\nabla} u_3)$$

vector identity $\vec{\nabla} \cdot (\phi \vec{B}) = \vec{\nabla} \phi \cdot \vec{B} + \phi \vec{\nabla} \cdot \vec{B}$

$\phi \rightarrow$ scalar

$$\phi = A_1 h_2 h_3 ; \vec{B} = \vec{\nabla} u_2 \times \vec{\nabla} u_3$$

And let us use this e_1 and let us substitute it here to expand the divergence. So, this will become the divergence of $A_1 e_1$ cap, this is equal to a divergence of $e_1, h_2 h_3$, grad u_2 cross

grad u_3 . And remember this identity is very, this we have, similar identities we can write for the other direction. So, we can write, let us say e_2 , e_2 is h_3 , h_1 , grad u_3 cross grad u_1 , just using the cyclic nature of 123 and similarly, for e_3 , we can write e_3 equals h_1 , h_2 , grad u_1 cross grad u_2 . So, let us see now. So, let us remember vector identity which we know from before, vector identity, which is basically the equivalent of kind of the product rule of calculus.

So, let us say we have the divergence of a scalar function let us say, ϕ times another vector let us say B . So, ϕ is a scalar and B is a vector. So, this becomes, remember this is like the product rule, so it is like the grad ϕ dot B plus ϕ divergence of B . So, if we use this identity for $A_1 e_1$ where we recognize that what is ϕ and what is B , so, if you see ϕ , like so, if we compare let us say ϕB , we identify that ϕ , if we represent ϕ by this term, which is $A_1 h_2 h_3$ and B vector by a cross product of grad u_2 cross grad u_3 .

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vector identity $\nabla \cdot (\phi \vec{B}) = \nabla \phi \cdot \vec{B} + \phi \nabla \cdot \vec{B}$
 $\phi \rightarrow$ scalar
 $\phi = A_1 h_2 h_3$; $\vec{B} = \nabla u_2 \times \nabla u_3$
 $\nabla \cdot (A_1 \hat{e}_1) = \nabla \cdot (A_1 h_2 h_3 \nabla u_2 \times \nabla u_3)$
 $= \nabla(A_1 h_2 h_3) \cdot (\nabla u_2 \times \nabla u_3) + A_1 h_2 h_3 \nabla \cdot (\nabla u_2 \times \nabla u_3)$
 $= \nabla(A_1 h_2 h_3) \cdot \frac{\hat{e}_2}{h_2} \times \frac{\hat{e}_3}{h_3} = \nabla(A_1 h_2 h_3) \cdot \frac{\hat{e}_1}{h_2 h_3}$
 $= \left[\frac{\hat{e}_1}{h_1} \frac{\partial}{\partial u_1} (A_1 h_2 h_3) + \frac{\hat{e}_2}{h_2} \frac{\partial}{\partial u_2} (A_1 h_2 h_3) + \frac{\hat{e}_3}{h_3} \frac{\partial}{\partial u_3} (A_1 h_2 h_3) \right] \cdot \frac{\hat{e}_1}{h_2 h_3}$

So, recognizing this is ϕ and B and using this vector identity we can rewrite the divergence of $A_1 e_1$, so, that is the divergence of $A_1 h_2 h_3$ grad u_2 cross grad u_3 , that is what we had before and now using the rule of this vector and identity, recognizing appropriately what is ϕ and what is B , we can expand this as the grad of $A_1 h_2 h_3$ dot product with grad u_2 cross grad u_3 plus $A_1 h_2 h_3$ divergence of vectors is grad u_2 cross grad u_3 .

And now realizing what this can be, so, this is like the divergence of curl or something you can think about, this is like. It is not strictly like that, but you can think about in this direction and we can show that this will be equal to 0. So, then this expression becomes grad of $u_1 h_2 h_3$ dot and grad of u_2 we can use this identity. This identity we can write as e_2 by h_2 that

can be written like e_2 by h_2 , using the contravariant components instead of the covariant. This will be written as e_3 by h_3 .

So, finally, this could be written as a grad of $A_1 h_2 h_3$ and e_2 cross e_3 , that is e_1 divided by $h_2 h_3$. So, just like this, we can write the other components basically. But, before writing that, let us expand this thing even further. So, remember the grad operator, the grad operator is defined using this, this del operator acting. So, this is the grad acting on a scalar functions, this is the sum over i equals 1 to 3 e_i by h_i .

So, using that we can expand this basically, we can expand that. And that becomes e_1 by h_1 partial u_1 , e_1 , $h_2 h_3$ and we have a sum, so three components e_2 by h_2 partial u_2 $A_1 h_2 h_3$ and then plus e_3 over h_3 partial u_3 $A_1 h_2 h_3$ and this entire thing dot product with $e_1 h_2 h_3$. So, if you see we are just dotting it with e_1 , that is the only component which will survive. So, this is the thing that will survive, others will become 0.

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$$\vec{\nabla} \cdot (A_1 \hat{e}_1) = \frac{1}{h_1 h_2 h_3} \frac{\partial}{\partial u_1} (A_1 h_2 h_3)$$

$$\vec{\nabla} \cdot (A_2 \hat{e}_2) = \frac{1}{h_1 h_2 h_3} \frac{\partial}{\partial u_2} (A_2 h_3 h_1)$$

$$\vec{\nabla} \cdot (A_3 \hat{e}_3) = \frac{1}{h_1 h_2 h_3} \frac{\partial}{\partial u_3} (A_3 h_1 h_2)$$

$$\vec{\nabla} \cdot \vec{A} = \frac{1}{h_1 h_2 h_3} \left[\frac{\partial}{\partial u_1} (A_1 h_2 h_3) + \frac{\partial}{\partial u_2} (A_2 h_3 h_1) + \frac{\partial}{\partial u_3} (A_3 h_1 h_2) \right]$$

So, the orthogonality of e vectors- $e_1 e_2 e_3$ conditions. Then this will become divergence, let us say divergence of $A_1 e_1$ cap, this is equal to you see, $h_2 h_3$ and h_1 is down. So, it will become 1 over $h_1 h_2 h_3$ and then this will become partial respect to u_1 and this will become $A_1 h_2 h_3$ as a result. Similar expressions we can have for let us say the other components, let us say $A_2 e_2$, I am just using the cyclic nature of things, we can write it as $h_1 h_2 h_3$ and this will be partial partial u_2 and this will become A_2 and you see the cycle 123.

So, this will have the cycle 231 $h_3 h_1$. And similarly, we can write divergence of $A_3 e_3$ which is 1 over $h_1 h_2 h_3$ and then this will become partial u_3 and this will become $A_3 e_3$ and

instead of, so that will become $h_1 h_2$. So, that is the divergence of $A_1 e_1$ cap $A_2 e_2$ cap and $A_3 e_3$ cap and now identifying like we can write the vector A in the contravariant basis like $A_1 e_1$ cap plus $A_2 e_2$ cap plus $A_3 e_3$ cap, we can take all the all the things together to find out the divergence of the vector field in an arbitrary curvilinear coordinates with scale factor $h_1 h_2 h_3$ as this.

So, this is $h_1 h_2 h_3$ and then we will have as a sum basically, so as you can see, it become partial $u_1 A_1 h_2 h_3$ plus partial $u_2 A_2 h_3 h_1$ plus partial $u_3 A_3 h_1 h_2$. So, this is the general expression for the divergence of a vector written in an arbitrary curvilinear coordinate system where we realize $h_1 h_2 h_3$ as the appropriate scale factors in the 123 direction respectively. So, this is the expression that we were targeting to have.

This is a very, very general expression as you will see and this expression we will see, use it for various coordinate systems to get the expressions for the divergence operator in the corresponding coordinate system.

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The image shows a handwritten derivation on a grid background. At the top, it states the divergence of a vector field $\vec{\nabla} \cdot \vec{A}$ in curvilinear coordinates as:

$$\vec{\nabla} \cdot \vec{A} = \frac{1}{h_1 h_2 h_3} \left[\frac{\partial}{\partial u_1} (A_1 h_2 h_3) + \frac{\partial}{\partial u_2} (A_2 h_3 h_1) + \frac{\partial}{\partial u_3} (A_3 h_1 h_2) \right]$$

Below this, it shows the substitution $\vec{A} = \vec{\nabla} \phi$ and the definition of the components $A_i = \frac{1}{h_i} \frac{\partial \phi}{\partial u_i}$. This leads to the expression for the Laplacian of a scalar field ϕ :

$$\vec{\nabla} \cdot \vec{\nabla} \phi = \nabla^2 \phi = \Delta \phi$$

$$\Delta \phi = \frac{1}{h_1 h_2 h_3} \left[\frac{\partial}{\partial u_1} \left(\frac{h_2 h_3}{h_1} \frac{\partial \phi}{\partial u_1} \right) + \frac{\partial}{\partial u_2} \left(\frac{h_3 h_1}{h_2} \frac{\partial \phi}{\partial u_2} \right) + \frac{\partial}{\partial u_3} \left(\frac{h_1 h_2}{h_3} \frac{\partial \phi}{\partial u_3} \right) \right]$$

So, now, if in this expression, let us say if we substitute instead of A , if we substitute grad of any scalar, that means, what we have in my mind is that EI is basically remember from partial ϕ partial U_i , the i th component of the... this is what we are talking about, grad of ϕ . So, that is the components.

So, if we recognize this, that A_i is this, this is the gradient of ϕ . So, this is the thing that we are thinking about, this expression basically, if you see here, instead of f , we have A_1 over h_1 . So, that is the idea that we are thinking. So, if we substitute A is equal to grad ϕ , you see,

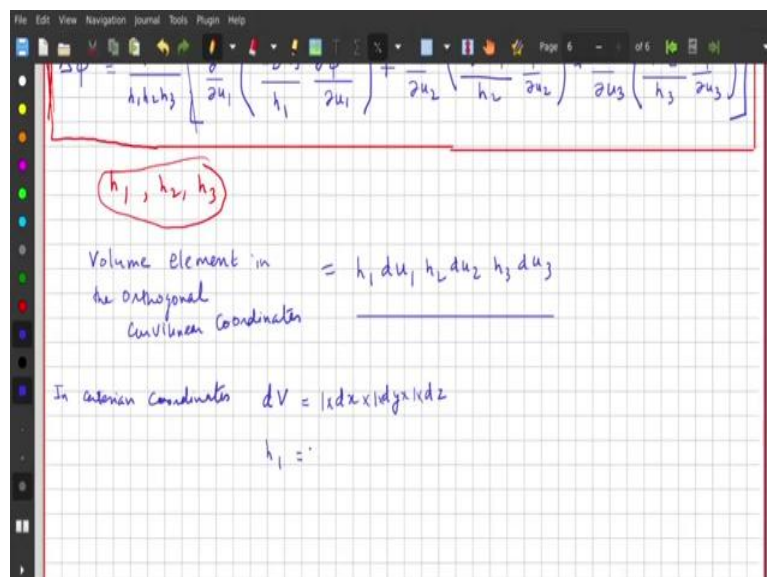
the left hand side becomes the divergence of the gradient of phi, so this becomes the Laplacian operator that is square phi or that can be written as this is a Laplacian of phi, so we can have a general expression for the Laplacian of any scalar field.

So, the general expression for the Laplacian of a scalar field is given by, just becomes 1 over $h_1 h_2 h_3$, it naturally follows from this expression which we had and this will become partial u_1 and instead of A_1 , we will have 1 over h_1 partial phi partial u_1 basically. So, we will have $h_2 h_3$ which are these two things, and then we have one h_1 from here, and then we will have partial phi partial u_1 . And like that for all the others, we just substitute as A_i s appropriately here.

So, this will become partial with respect to u_2 , and this is $h_3 h_1$ by h_2 partial phi partial u_2 and plus partial with respect to u_3 . And this is $h_1 h_2$ by h_3 partial phi partial u_3 . So, that is what we have as a general expression for the Laplacian. So, we have written the general expression for the divergence and we have written the general expression for the Laplacian of a scalar field.

So, the divergence has been a vector and the Laplacian of a scalar field in arbitrary curvilinear coordinate. So, these two expressions are very powerful. And as you see, the only thing that we require to get to a specific form is knowing what are the scale factors.

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So, if we know what are h_1 , h_2 and h_3 , we can get the specific expression for the individual coordinate system. That is the power of this, this method and the generalized that we are

trying to build. So, now we have these two expressions and using these we can write the divergence of a vector field and the Laplacian of a vector field in arbitrary coordinate systems.

So, we can write the divergence and the Laplacian in a particular coordinate system, this specific coordinate system and that is what we will be doing now in seeing the specific cases. So, before that, let us write what is the volume element in the orthogonal curvilinear coordinates that we used? And that is $h_1 du_1 h_2 du_2 h_3 du_3$. That is the volume element. So, if we compare this with the volume element in each coordinate system, which we had before, we will be able to figure out what h_1, h_2 and h_3 are in general.

And another thing to, why these two expressions are so powerful? It looks very long, but these are very, very simple, there is a pattern to it. If you see like h_1, h_2, h_3 and then partial with respect to u_1, u_2 and then 123, 231, 312, the cyclic nature is there. So, that makes it very easy to like keep this thing in mind. So, let us say for example, in Cartesian coordinates, we realized what was the volume elements, the volume element dV was $dx dy dz$. And if you recall this, we can rewrite like this, where we realize the individual $h_1 h_2 h_3$, the scale factors all are one.

So, h_1, h_2, h_3 all are one and just if you substitute h_1, h_2, h_3 all equal to 1, you will see divergence of the vector as well as the Laplacian, you will get the familiar fields, the familiar Cartesian expressions.

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The image shows handwritten notes on a grid background. The text is as follows:

In cartesian coordinates $dV = |x| dx |y| dy |z| dz$
 $h_1 = h_2 = h_3 = 1$

In cylindrical coordinates $dV = (dr) (r d\theta) dz$
 $h_1 = 1$ ($\because h_1 du_1 = 1 dr$)
 $h_2 = r$
 $h_3 = 1$

$\Delta \phi = \frac{1}{x} \left[\frac{\partial}{\partial x} \left(x \frac{\partial \phi}{\partial x} \right) + \frac{\partial}{\partial \theta} \left(\frac{1}{r} \frac{\partial \phi}{\partial \theta} \right) + \frac{\partial}{\partial z} \left(r \frac{\partial \phi}{\partial z} \right) \right]$

$\Delta \phi = \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial \phi}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 \phi}{\partial \theta^2} + \frac{\partial^2 \phi}{\partial z^2}$

So, let us do another case where let us say in cylindrical coordinates, what is the volume element? Let us write the volume elements. So, dV in cylindrical coordinate is dr into $r dr$

or no, if you remember this is dr into $r d\theta$. This was dr into $r d\theta$ that was in the azimuthal direction and times dZ . And if we compare this now with this expression, we will see that h_1 is 1 since what we are doing is $h_1 du_1$ is basically dr where we recognize e_1 is R and the factor that sits in front of dr is 1.

So, h_1 is 1. Similarly, we will get h_1 is 1, h_2 is r and h_3 is 1. So, if we use the values of $h_1 h_2 h_3$ in these two expressions, we will get the appropriate expression for the divergence as well as the Laplacian. Let me write the Laplacian operator in cylindrical coordinates. So, that will become 1 over r , because remember you have 1 over $h_1 h_2 h_3$ and you see h_1 and h_3 is 1 and h_2 is r . So, the product is r and this becomes $\nabla \cdot \nabla r$ of $r \nabla \phi \nabla r$. You see, because that is $h_2 h_3$ by h_1 and this h_2 is r , this r comes from that h_2 plus $\nabla \cdot \nabla \theta$ 1 over $r \nabla \phi \nabla \theta$ plus $\nabla \cdot \nabla Z$ $r \nabla \phi \nabla Z$.

And if we know, just open the brackets we will have Laplacian of the scalar field, ϕ in cylindrical coordinates is given by 1 over $r \nabla \cdot \nabla r$ $r \partial^2 \phi \partial r^2$ plus 1 over r^2 into $\phi \partial^2 \theta$ plus $\Delta \phi \partial Z^2$, so that is the Laplacian in cylindrical coordinate systems. Similarly, we can write it for spherical polar coordinates. So, let us do that.

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Spherical Polar Coordinates $dV = (dr)(r d\theta)(r \sin\theta d\phi)$

$h_1 du_1 \quad h_2 du_2 \quad h_3 du_3$

$h_1 = 1; h_2 = r; h_3 = r \sin\theta$

$$\Delta \phi = \frac{1}{r^2 \sin\theta} \left[\frac{\partial}{\partial r} \left(r^2 \sin\theta \frac{\partial \phi}{\partial r} \right) + \frac{\partial}{\partial \theta} \left(\sin\theta \frac{\partial \phi}{\partial \theta} \right) + \frac{\partial}{\partial \phi} \left(\frac{1}{\sin\theta} \frac{\partial \phi}{\partial \phi} \right) \right]$$

$$\Delta \phi = \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial \phi}{\partial r} \right) + \frac{1}{r^2 \sin\theta} \frac{\partial}{\partial \theta} \left(\sin\theta \frac{\partial \phi}{\partial \theta} \right) + \frac{1}{r^2 \sin^2\theta} \frac{\partial^2 \phi}{\partial \phi^2}$$

So, let us write it for spherical polar coordinates. Let us identify what was the volume element. So, if you recall the volume element dV was dr into $r d\theta$ into $r \sin\theta d\phi$, now this was the azimuthal, this was the zenith term and this was the radial term. And if with this we recognize with $h_1 du_1 h_2 du_2 h_3 du_3$, we realize that h_1 is 1, h_2 is r and h_3 is $r \sin\theta$.

theta. And if you use this $h_1 h_2 h_3$ you can show that this will be, just substitute $h_1 h_2 h_3$ in this equation and out pops the Laplacian of a scalar field in spherical polar coordinates.

Because of this will become $\frac{1}{r^2 \sin \theta} \frac{\partial}{\partial r} (r^2 \sin \theta \frac{\partial \phi}{\partial r}) + \frac{\partial}{\partial \theta} (\frac{1}{r \sin \theta} \frac{\partial \phi}{\partial \theta}) + \frac{1}{r \sin^2 \theta} \frac{\partial^2 \phi}{\partial \phi^2}$. That is laplacian in spherical polar coordinates. We saw how to get each of the Laplacian operators in cylindrical as well as spherical polar using the general expression that we derived and same thing can be done for the divergence.

This grad of phi is $\frac{1}{r^2} \frac{\partial \phi}{\partial r} + \frac{1}{r^2 \sin \theta} \frac{\partial \phi}{\partial \theta} + \frac{1}{r^2 \sin^2 \theta} \frac{\partial \phi}{\partial \phi}$. That is laplacian in spherical polar coordinates. We saw how to get each of the Laplacian operators in cylindrical as well as spherical polar using the general expression that we derived and same thing can be done for the divergence.

So, that is what we wanted to show, a general way of writing these two expressions are the things that we wanted to understand and then we wanted to show how general these two expressions are for the divergence as well as the Laplacian of a vector and a scalar respectively in arbitrary curvilinear coordinate systems. So, that is for this segment and we will see you in the next segment with some other materials. So, thank you.