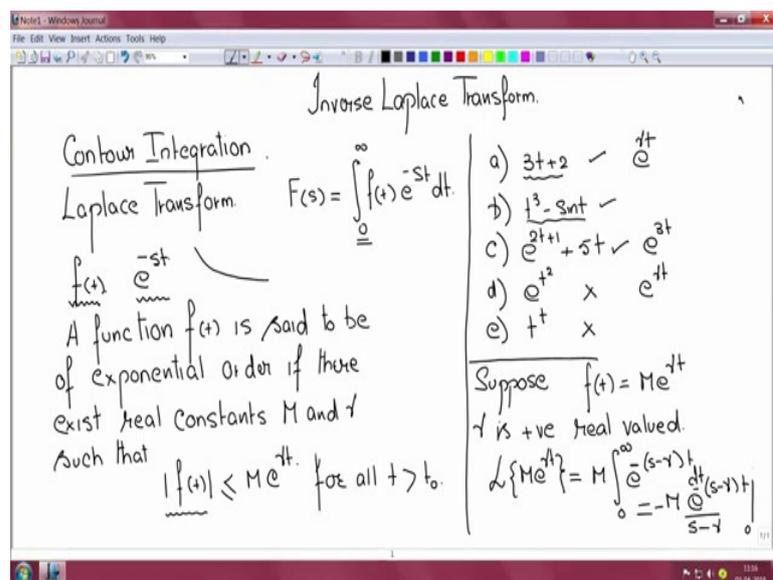


A short lecture series on Contour Integration in the Complex Plane
Prof. Venkata Sonti
Department of Mechanical Engineering
Indian Institute of Science, Bengaluru

Lecture - 22
Inverse Laplace Transform

Good morning, welcome to this lecture on complex variables. So, last class we looked at the L shaped branch cut, that was meant as at a research level not at a classroom level, but still it is very rarely discussed or presented, so I have given hints as to how an L shaped cut can come. Now, today I discuss the last topic in this series which is Inverse Laplace Transform.

(Refer Slide Time: 00:59)



Typically, Laplace transforms are used in engineering in solving ODEs and PDEs and the inverse Laplace transform is used from a set of tables. So, we rarely actually compute the inverse Laplace transform, but it turns out that inverse Laplace transform requires contour integration.

So, I will show you how to get an inverse Laplace transform using contour integration. So, just to begin Laplace transforms, the forward Laplace transform is typically given by capital F, s is the Laplace symbol variable, integral 0 to infinity f of t which is the time signal, e to the power of minus s t d t ok. Our systems are causal, they are excited at time t equal to 0 and so this limit is 0 and it goes off to infinity, ok.

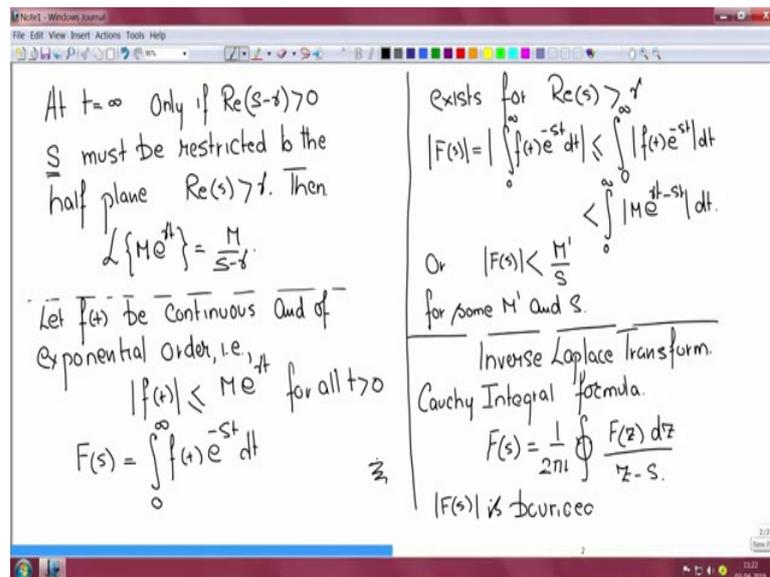
Now, in a simplistic sense the Laplace transform will exist if the function f of t does not grow faster than e to the power minus $s t$ can bring it down ok, e to the power minus $s t$ is a decaying type function. So, if f of t does not grow faster, so much faster that e to the power minus $s t$ cannot pull it down, such functions have Laplace transform, and in engineering it is not a very big issue to find such such responses, ok. All our response is finally decay to 0.

Now, some formal presentation; a function, a function f of t is said to be of exponential order if there exist real constants M and γ such that f of t amplitude is bounded in this manner; $M e$ to the power γt ok, e to the power γt is an increasing function, rapidly increasing function. So, f of t amplitude is bounded in this fashion, ok.

For all, so we will write for all t greater than some t_0 . So, in this sense if you look at the following functions a), $3 t$ plus 2 and lets say b) is t cube minus $\sin t$; c) e to the power $2 t$ plus 1 plus $5 t$, then d) e to the power t square and e) t to the power t . If you look at these functions these three are of exponential order. It is very easy to find an e to the power t or some γt which bounds this function and this function, and of course, this function is easily bounded by $3 t$, but these two are not of exponential order, they rise much faster than any e to the power γt .

Suppose, suppose, f of t is indeed equal to $M e$ to the power γt , where γ is positive real valued, ok, then the Laplace transform of $M e$ to the power γt is equal to $M \int_0$ to infinity, e to the power minus s minus γt dt which is equal to, minus $M e$ to the power minus s minus γt by s minus γ over the limits 0 to infinity.

(Refer Slide Time: 08:13)



Now, at t equal to infinity, the expression will be bounded; the above expression is bounded only if the real part of s minus γ is positive, ok, which means s the Laplace variable is must be restricted to the half plane, real part of s greater than γ . Then Laplace transform of $M e^{\gamma t}$ is equal to M by s minus γ , ok.

Further let $f(t)$ be continuous and of exponential order, that is magnitude of $f(t)$ is less than or equal to $M e^{\gamma t}$, for some M and some γ , let us say for all $t > 0$. Then the Laplace transform $F(s)$ is equal to, integral 0 to infinity $f(t) e^{-st} dt$ exists for real part of s greater than γ .

And we see magnitude of $F(s)$ is equal to the magnitude of the integral 0 to infinity, $f(t) e^{-st} dt$, it is less than equal to the integral 0 to infinity, the magnitude of $f(t) e^{-st} dt$, which is less than strictly integral 0 to infinity $M e^{\gamma t - st} dt$. Or magnitude of $F(s)$ is definitely less than M dash by s , for some M dash and some s .

Now the inverse Laplace transform, inverse Laplace transform. We use the Cauchy integral formula to write it, Cauchy integral formula, we write it as $F(s)$ is equal to 1 over twice πi , integral a counterclockwise contour $F(z) dz$, by z minus s , ok.

(Refer Slide Time: 14:45)

$|F(s)|$ is Bounded & Analytic
 for $\text{Re}(s) > \gamma$

$F(s) = \frac{1}{2\pi i} \oint_{CC} \frac{F(z) dz}{z-s}$

$F(s) = \frac{1}{2\pi i} \oint_{CC} \frac{F(z) dz}{s-z}$

$F(z)$ is Laplace Transform decays as $\frac{M}{|z|}$ and as $|z| \rightarrow R \rightarrow \infty$

$\frac{F(z) dz}{s-z} \sim \frac{M}{R e^{i\theta}} \cdot \frac{1}{R e^{i\theta}} \cdot R e^{i\theta} d\theta$

$\left| \frac{F(z) dz}{s-z} \right| \rightarrow \left| \frac{M}{R} \cdot \frac{1}{R} d\theta \right| \rightarrow 0$ as $R \rightarrow \infty$

we are left with

$F(s) = \frac{1}{2\pi i} \int_{C+100} \frac{F(z) dz}{s-z}$

Let us remember that F of s is bounded, is bounded, is bounded; let me change the page, F of s is bounded and is analytic for real values of s greater than γ . Now let us use, so we are going to look at that integral and let us use the following contour. The contour is this, we have the complex plane and we have γ over here, here is my γ . So, I take a contour counterclockwise, so I come down from infinity above to infinity below, I go around this infinite semicircle, ok.

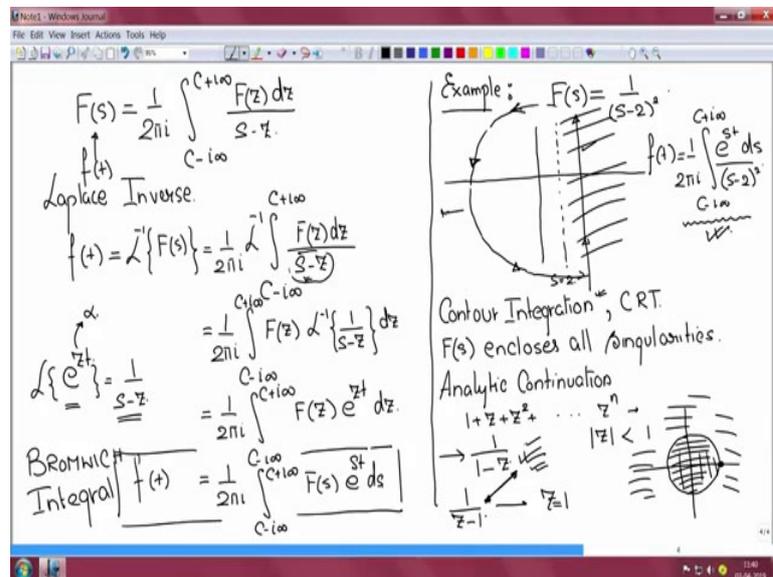
So, let me remind you that what we wrote was F of s is equal to 1 over twice πi integral counterclockwise F of z dz , by z minus s , ok. Now we switch it to clockwise and change the sign. So, F of s is equal to 1 over twice πi , we switch it to clockwise F of z dz over s minus z . So, you brought in two negatives, we changed the direction and we change the sign on here. Now, F of z is the Laplace transform and decays as bounded by M dash over magnitude of z as before and z tends to infinity over a radius R the integrand decays.

So, integrand here is F of z , by s minus z , ok. F of z is bounded by M dash over $R e$ to the power of i theta. Then s minus z , s is less important as z tends to infinity. So, I have $R e$ to the power of i theta and then dz is $R i e$ to the power of i theta $d\theta$. So, we are essentially looking at the integrand over this semicircle. So, z is parametrized at $R e$ to the power i theta. So, now, the magnitude of this F of z neither should be here, F of z dz

by s minus z behaves as M over R , 1 by R and $R d\theta$ which tends to 0 as R tends to infinity.

So, on this contour, this portion of the integral will not survive, on this contour here, on this contour here, this portion of the integral will not survive; that is what we have shown. On this semi circular contour, z is equal to $R e^{i\theta}$, so this integral integrand goes to 0 . So, now we are left with the straight line portion so, which is what F of s is equal to 1 over twice πi . We are to the right of γ over here, this is a straight line so, we will write that as some C minus i infinity to C plus i infinity F of z dz over s minus z . I will write this again here.

(Refer Slide Time: 21:23)



So, what we have is F of s is equal to 1 over twice πi , integral the straight line portion C minus i infinity to C plus i infinity, F of z dz by s minus z . Now, let us take the Laplace inverse of this, because it has arisen from time function, ok. So, now the f of t which is Laplace inverse of F of s is equal to 1 over twice πi , the Laplace inverse C minus i infinity to C plus i infinity, F of z dz by s minus z , ok.

Now, Laplace will inverse will act only on s and not on z . So, what we have here, is 1 over twice πi , integral C minus i infinity to C plus i infinity, F of z is unaffected, then the Laplace inverse goes and operates on 1 over s minus z .

We can compute the basic Laplace transform of this function e^{-zt} and the Laplace transform of this is $1/(s-z)$, ok. Normally we are used to seeing α over here if you remember your Laplace transforms. So, now the Laplace inverse of this is e^{-zt} . So, we get $1/(2\pi i) \int_{C-i\infty}^{C+i\infty} F(z) e^{-zt} dz$ or using another dummy variable, $1/(2\pi i) \int_{C-i\infty}^{C+i\infty} F(s) e^{-st} ds$.

So, that is how the Laplace inverse is computed actually over a straight line segment where $F(s)$ is analytic. The straight line segment going from $-i\infty$ to $+i\infty$ and this is called the Bromwich integral. Now, let us see; now let us consider take an example, the example is this suppose I have, so let's take an example, the example is this; $F(s)$ is equal to $1/(s-2)^2$.

Those of you who remember your Laplace tables may immediately know what $f(t)$ is. So, now, here I have a pole at $s=2$. So, now we said that $F(s)$ is analytic in the region as s is greater than some γ , so you can see that $s=2$ will be our γ here, this is $s=2$ and so now, we have a straight line segment that goes from this way upwards. So, now we have to compute this integral; $f(t)$ is equal to $1/(2\pi i) \int_{C-i\infty}^{C+i\infty} F(s) e^{-st} ds$, over $s-2$ whole squared, ok.

Now, using contour integration; we have to use contour integration that is the whole thesis of this series of lectures. We are going to use contour integration and again as before, this segment which we want we do not touch, so we will use principles of contour integration and if possible Cauchy residue theorem and compute this integral without touching it, ok. So, now what we do, we close this, we close this in the left half plane of the; we close it in the left half complex plane so, we now have a closed contour.

So, that now the function $F(s)$ encloses all its singularities. Now one important point here which I will briefly mention is that in the beginning we said that $F(s)$ exists and is analytic to the right of a certain region. So, we should ideally sit only in that region and not extend it to the left, but here is an idea called analytic continuation which is used. I will explain it using a simple example. Suppose we have this geometric series $1+z+z^2+\dots$ and so forth and goes off to infinity; the sum we know is $1/(1-z)$ provided that magnitude of z is less than 1.

That means what in the complex plane, you have a unit circle, you have a unit circle, and so you have to be inside the unit circle, so that you have convergence, ok; however, separate from this idea where we sum this geometric series, there exists a function $1/(z-1)$ or $1/(1-z)$ which is singular only at $z=1$ and is well behaved, everywhere else, there exists a function like that, ok.

Now, the idea of analytic continuation says that this $1/(1-z)$ and this $1/(1-z)$ are identical. They are identical, they are not two separate functions, so that is the idea of analytic continuation. So, now this function which is supposed to converge inside the $|z|=1$ radius extends and is singular only at $z=1$, ok. So, that is the idea of analytic continuation. Now we are going to do this contour integral and as before make sure that the function or the integrand goes to 0 on this arc, we will make sure of that and use the residue theorem. Right now time is running out, so we will continue in the next class.

Thank you.