

# APPLIED ELASTICITY

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WEEK: 03

Lecture- 11

COURSE ON:  
APPLIED ELASTICITY

Lecture 11  
STRESS MEASURES I

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Welcome back to the course of applied elasticity. In the third week of this lecture, we are going to first discuss about the different stress measures. In the previous week, we had talked about the strain and displacement measures kinematics of deformation. Now, we are moving forward to the stress measures within a body.

Surface Traction

Surface traction is defined as force per unit area acting on a plane with unit normal  $\hat{n}$  at any point  $P$ .

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First, we are going to define the surface traction and based on that the stress components would be defined.

So, surface traction is defined to be the force acting per unit area on a plane at any point  $P$  and  $\hat{n}$  being the unit normal to the plane. So, for defining surface traction, we need a plane defined by its unit normal direction  $n$  vector and a point on the plane. If you are just given with a single point without any plane, it is not possible to define the surface traction and thus it is not possible to define the stress components as well. So, for defining surface traction or stress component we need a point and the plane on which we are going to define the surface traction vectors.

Let us consider a body which is supported by different boundary conditions, pin support, roller support, rigid support all different types of supports may be there within this continuum. And this is subjected to different types of external loadings  $F_1, F_2, F_3$  are point loads or concentrated loads acting at different points and small  $f$  is a distributed load acting for a specific region. Now,  $P$  is one point of interest within the body and we are interested to find out the surface traction at point  $P$ .

By its definition, for writing the surface traction expressions or defining it, we first need to define a plane. So, let us cut this body using a plane as shown by this red curve which passes through point  $P$ . and this divides the body into two halves part  $A$  and part  $B$ . So, using this particular plane passing through point  $P$  we are dividing the body into part  $A$  and part  $B$  and then we are going to draw the free body diagram of both the parts separately and this plane is defined by its unit outward normal vector small  $\hat{n}$ .

Now, first starting with the free body diagram of part  $A$  The supports and forces are shown whatever were attached to part  $A$ . Now, on this cutting plane, we are considering a small elemental area  $\Delta A$  around that point  $P$  and  $\Delta A$  is the small elemental area. This plane, the open surface of this plane is denoted by surface  $\Gamma$ , capital  $\Gamma$ .  $\hat{n}$  is the unit normal vector to this surface  $\Gamma$  at point  $P$  and let us say  $f \Delta \tilde{F}$  this is the elementary force acting at this particular point due to cutting of the body with this particular plane defined through  $n$  vector.

Now, similarly on the other part, part  $B$ , if you draw the free body diagram, it would be looking like this. There will be an identical  $P$  point on part  $B$  around which we are taking elemental area  $\Delta A$  or  $\Delta A$  and that is having unit outward normal of  $-\hat{z}$ . Opposite to the outward normal of part  $A$  surface, we will be having equal magnitude, but opposite direction unit outward normal vector on the open surface or cutting plane of part  $B$  of the body. On the point  $P$  of part  $B$  minus of  $\Delta \tilde{F}$  amount of force is acting which is negative

of the same amplitude force as it was acting on the part  $A$  point  $P$  on the small elemental area  $\Delta A$ .

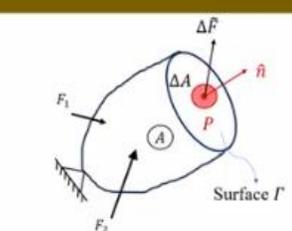
**Surface Traction**

The traction vector  $\tilde{t}^{(n)}$  acting on a plane with unit normal  $\hat{n}$  at point  $P$  is defined as

$$\tilde{t}^{(n)} = \lim_{\Delta A \rightarrow 0} \frac{\Delta \vec{F}}{\Delta A}$$

Force exerted by region  $B$  on region  $A = \int_{\Gamma} \tilde{t}^{(n)} dA$

If  $\tilde{t}^{(-n)}$  is the traction vector on opposite side of the surface on region  $B$ , then that is equal and opposite to  $\tilde{t}^{(n)}$  as,

$$\tilde{t}^{(n)} = -\tilde{t}^{(-n)}$$


Components of  $\tilde{t}^{(n)}$  along  $\hat{n} \rightarrow$  Normal stress  
 Components of  $\tilde{t}^{(n)}$  along  $\Delta A \rightarrow$  Shear stress

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Now coming to the definition of the surface traction that surface traction vector  $\tilde{t}^{(n)}$  this superscript  $n$  denotes the normal normal to the plane on which this surface traction is defined so  $\tilde{t}^{(n)}$  is acting on a plane with unit normal  $\hat{n}$  at point  $P$  and that is defined as limit  $\Delta A$  tending to 0  $\Delta \vec{F}$  by  $\Delta A$ . So, elementally elemental force acting at point  $P$  divided by elemental area around point  $P$  with limit this elemental area  $\Delta A$  tending to 0 is defined to be the surface traction  $\tilde{t}$  on this particular plane defined by unit normal  $\hat{n}$ . if you take the components of this surface traction vector  $\tilde{t}^{(n)}$ . So, surface traction or  $\tilde{t}^{(n)}$  is a vector quantity because this is limit  $\Delta A \Delta \vec{F}$  vector by  $\Delta A \Delta \vec{F}$  being a force vector that force vector per unit area is defined to be surface traction vector and thus this is a vector quantity. So, component of surface traction vector along the normal direction is called the normal stress. Component of the surface traction vector along the surface  $\Gamma$  parallel to this area  $\Delta A$  that is called the shear stress. So, total surface traction can be can be obtained by taking the taking the resultant of the normal component and the shear component. Now, if you integrate this surface fraction over the entire surface  $\Gamma$ , so  $\tilde{t}^{(n)} dA$  integrated over entire  $\Gamma$  gives us the total force exerted on surface  $\Gamma$  of part  $A$  due to part  $B$  or region  $B$ . So, this much amount of force the region  $B$  exerts on region  $A$  part  $A$  once we are cutting it with this plane with unit normal  $n$  that total force is integral  $\Gamma \tilde{t}^{(n)} dA$ .

Now,  $\tilde{t}^{(-n)}$  being the traction vector on opposite side of the surface of part  $B$ , that should be equal and opposite to the surface traction acting on part  $A$ . So, as already discussed  $\tilde{t}^{(n)}$  will be minus of  $\tilde{t}^{(-n)}$ , the magnitude wise these two surface tractions are identical, but their directions would be just opposite and thus Once we are combining these two cut

sections part  $A$  and part  $B$  together once again this  $\tilde{t}^{(n)}$  would cancel minus of  $\tilde{t}^{(-n)}$  and thus in the total body there is no force no traction acting at particular internal point  $P$ . Once you are cutting it then only that will become open surface and for both part  $A$  and part  $B$  equal and opposite surface traction forces would be acting.

**Cauchy Stress**

FBD of an infinitesimal cube centered at  $P$  with edges parallel to the coordinate axis

- This is a stress measure in deformed configuration.
- For an infinitesimal cube centered at point  $P$ , the Cauchy stress components are defined as the tractions acting on the external surfaces of this cube.
- The components of Cauchy stress  $\sigma_{ij}$  is the traction acting on a face with normal along  $x_i$  direction and the component is along  $x_j$  direction.

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This particular equation is called Cauchy's lemma. Now, moving forward to the definition of Cauchy stress which is based on the definition of surface fractions we are considering a small infinitesimal cube and this Cauchy stress is the stress measure in the deformed configuration. So, we will be looking into the stress measures in undeformed configuration later. First, we are discussing the Cauchy stress components which are the stress components defined in the deformed configuration.

For this small cube, infinitesimal cube centered at point  $P$ , the Cauchy stress components are defined as the surface tractions acting on all six external surfaces of the cube. Now, Cauchy stress tensor or its component is denoted with two subscripts  $\sigma_{ij}$  and thus stress is a tensor quantity. As two subscripts are required to define stress, we must define this as a tensor quantity. We cannot use vector quantity for defining stress. As long as it is traction, that is a vector.

Using traction, once we are going to define stress, it must be a second order tensor. Now, what is the significance of these two subscript  $i$  and  $j$  in any Cauchy stress tensor component  $\sigma_{ij}$ ? Now, two subscript out of which the first one is  $i$ . which refers to the traction acting on a phase with normal along  $x_i$ . So, along  $x_i$  direction if the normal is aligned that defines the plane on which this Cauchy stress component  $\sigma_{ij}$  is defined and second subscript  $j$  refers to the direction.

So, first subscript refers to the plane whose normal is  $x_i$ , second subscript refer to the direction of the Cauchy stress component which is along  $x_j$  direction. Now, let us draw the Cauchy stress components for this cube. So, considering  $x_1$ ,  $x_2$  and  $x_3$  direction, note that I have used three different colours for three directions  $x_1$ ,  $x_2$ ,  $x_3$  which are deformed coordinates because Cauchy stress is defined in the deformed coordinate frame.

Now, considering the  $x_1$  planes, So, if you consider this particular plane this is having unit normal along the  $x_1$  direction. If you consider the back plane the left hand side plane of this cube that is also  $x_1$  plane its normal is along the minus  $x_1$  direction. For these two planes we are defining the Cauchy stress components like this. On the right hand side plane

As this plane is defined as  $x_1$  plane, the first subscript would be 1 for all the components because  $i$ , the first subscripts refer to the plane. For this particular plane, we are having it to be a  $x_1$  plane with unit normal along  $x_1$ . So, first subscript must be 1. So, for all these, first subscript is equals to 1. Now, second subscript refers to the direction.

This component is along one direction. So, that is why we are having  $\sigma_{11}$ . This component is along  $x_2$  direction. That is why name is  $\sigma_{12}$  and third component is along  $x_3$  direction. That is why name is  $\sigma_{13}$ .

On the negative  $x_1$  plane that is on the back side, this is negative  $x_1$  plane. Similarly,  $\sigma_{11}$ ,  $\sigma_{12}$ ,  $\sigma_{13}$  can be drawn. The directions are simply negative or opposite to the directions of the corresponding stress components on the positive  $x_1$  plane. So, like that you can define 3 stress components on each plane. For  $x_1$  plane, we are defining 3 stress components  $\sigma_{11}$ ,  $\sigma_{12}$ ,  $\sigma_{13}$  along  $x_1$ ,  $x_2$ ,  $x_3$  directions acting on  $x_1$  plane.

Similarly, we can move forward to rest of the planes. So, if you take the  $x_2$  planes which are denoted by these green colours. So, that is front face and back face of the cube. The first subscript is 2 because these are  $x_2$  plane. So, three stress components defined are  $\sigma_{21}$ ,  $\sigma_{22}$ ,  $\sigma_{23}$ .

The back face is the positive subscript.  $x_2$  phase. This is positive  $x_2$  plane whereas, the front is negative  $x_2$  plane. So, for the positive  $x_2$  plane,  $\sigma_{21}$ ,  $\sigma_{22}$ ,  $\sigma_{23}$  are along the positive  $x_1$ , positive  $x_2$ , positive  $x_3$  direction.

For the front face that is negative  $x_2$  plane, we are having  $\sigma_{21}$ ,  $\sigma_{22}$ ,  $\sigma_{23}$  along the negative  $x_1$ ,  $x_2$ ,  $x_3$  directions. Now, coming to the  $x_3$  plane that is the top plane and the bottom plane denoted by the blue colour stress components here the components are  $\sigma_{31}$   $\sigma_{32}$  and

$\sigma_{33}$  on both top and bottom positive  $x_3$  and negative  $x_3$  plane we can define the stress components. So, for  $\sigma_{ij}$  the first  $i$  refers to the direct to the plane with normal  $x_i$  second subscript  $j$  refers to the direction along  $x_j$ .

So, on the positive planes positive plane means planes whose normal is aligned along the positive  $x_i$  direction those are called positive planes. for positive planes along the positive direction  $\sigma_{ij}$  the cauchy stress component is taken to be positive On the negative planes along the negative direction we we are also taking  $\sigma_{ij}$  scotty stress components to be positive now with this we will move forward.

**Cauchy Stress Components**

$\vec{t}_1 = -\sigma_{11}\vec{e}_1 - \sigma_{12}\vec{e}_2 - \sigma_{13}\vec{e}_3$   
 $\vec{t}_2 = -\sigma_{21}\vec{e}_1 - \sigma_{22}\vec{e}_2 - \sigma_{23}\vec{e}_3$   
 $\vec{t}_3 = -\sigma_{31}\vec{e}_1 - \sigma_{32}\vec{e}_2 - \sigma_{33}\vec{e}_3$

FBD of an infinitesimal tetrahedron at P

$A_{PAB} = \frac{1}{2} dl_1 dl_3, A_{PAC} = \frac{1}{2} dl_1 dl_2, A_{PBC} = \frac{1}{2} dl_2 dl_3$   
 $\vec{AB} \times \vec{BC} = -dl_2 dl_3 \vec{e}_1 - dl_1 dl_3 \vec{e}_2 - dl_1 dl_2 \vec{e}_3$   
 $= -2 A_{ABC} \cdot \vec{n}$   
 Area of ABC =  $A_{ABC} = \frac{1}{2} |\vec{AB} \times \vec{BC}|$   
 $= \frac{1}{2} \sqrt{(dl_1, dl_2)^2 + (dl_2, dl_3)^2 + (dl_3, dl_1)^2}$

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Now to relate the traction vector with the corresponding Cauchy stress components, we are going to draw the Cauchy tetrahedron. So, this is a tetrahedron, where 3 sides are along  $x_1, x_2, x_3$  directions with unit vector  $\vec{e}_1, \vec{e}_2, \vec{e}_3$ . The tetrahedron is named with  $PABC$ . The side lengths  $PA$  equals to  $dl_1$ ,  $PB$  equals to  $dl_3$ ,  $PC$  equals to  $dl_2$ .  $dl_1, dl_2, dl_3$  are aligned along  $\vec{e}_1, \vec{e}_2, \vec{e}_3$  direction. Now, if you consider the inclined face  $ABC$ , That plane is inclined through all three axes whereas  $PAB$  plane is lying on  $x_1, x_3$  plane,  $PAC$  is lying on  $x_1, x_2$  plane,  $PBC$  is lying on  $x_2, x_3$  plane whereas  $ABC$  is one of the inclined plane.

Now, this for small infinitesimal tetrahedron defined as Cauchy tetrahedron around point  $P$ . Let us take  $A_i$  are the acceleration components along  $x_i$  direction,  $b_i$  are the body force component per unit mass along the  $x_i$  direction,  $\vec{t}^{(n)}$  is the traction vector acting on the inclined surface  $ABC$ . So, on the inclined surface, which is having  $n$  to be its unit normal vector,  $\vec{t}^{(n)}$  is defined that is surface traction on this inclined surface  $ABC$  defined by unit normal  $n$  and components of this  $\vec{t}^{(n)}$  traction vector are  $\tilde{t}_1, \tilde{t}_2, \tilde{t}_3, \tilde{t}_n$  vector becomes  $\tilde{t}_1 \vec{e}_1$  vector plus  $\tilde{t}_2 \vec{e}_2$  vector plus  $\tilde{t}_3 \vec{e}_3$  vector.

The traction components on other three phases  $PBC$ ,  $PAB$ ,  $PAC$ .  $PBC$  is nothing but  $x_1$  plane,  $PAB$  is nothing but  $x_2$  plane and  $PAC$  is  $x_3$  plane. On all those three planes, the traction components are defined with the help of corresponding Cauchy stress tensor components defined on those particular planes. So,  $\tilde{t}_1$  vector is that is the fraction component on  $PBC$  plane is equals to  $-\sigma_{11}\tilde{e}_1 - \sigma_{12}\tilde{e}_2 - \sigma_{13}\tilde{e}_3$ .

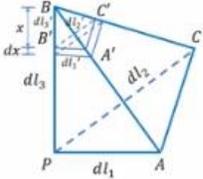
Why minus? Because all these components are acting on the negative  $x_1, x_2, x_3$  directions. Similarly, on plane  $PBC$ , you can define this traction vector  $\tilde{t}_2$ . On  $PAC$ , the bottom phase, you can define the traction vector  $\tilde{t}_3$  with the help of corresponding stress components like this. And this completes the overall free body diagram of the small tetrahedron around point  $P$ .

Now, if we are writing the expression of area of all these different sides, if you consider  $PAB$ , the front face, front vertical face in  $x_1, x_3$  plane, that is basically a triangle with base length  $dl_1$  and with height  $dl_3$ . of this particular triangle  $PAB$  is equals to half of base  $dl_1$  times  $dl_3$  that is written here. Area of  $PAB$  equals to half  $dl_1 dl_3$ . Similarly, area of  $PAC$  can be obtained as half of  $dl_1 dl_2$ , area of  $PBC$  can be obtained as half of  $dl_2 dl_3$ .

Moving forward if you take the cross product of  $\overrightarrow{AB}$  vector and  $\overrightarrow{BC}$  vector that can be written as minus of 2 times  $A_{ABC} \cdot \tilde{n}$  vector where  $\tilde{n}$  is the unit normal on the inclined plane  $\tilde{n}$ . So, by writing  $\overrightarrow{AB}$  vector and  $\overrightarrow{BC}$  vector as per its components as per their components and then taking the cross product you can show that  $\overrightarrow{AB}$  cross  $\overrightarrow{BC}$  equals to minus of two times area of inclined face  $ABC$  dot  $\tilde{n}$  vector. now from this  $A_{ABC}$  can be obtained as half of modulus of this cross product  $\overrightarrow{AB}$  and  $\overrightarrow{BC}$

so substituting the components of  $\overrightarrow{AB}$  and  $\overrightarrow{BC}$  from this equation the One component, first component of this cross product is  $dl_2, dl_3$ . Second component is  $dl_1, dl_3$ . Third component is  $dl_1, dl_2$ . Substituting all those here, the  $A_{ABC}$  inclined face can be obtained as half of square root of  $dl_1, dl_2$  whole square plus  $dl_2, dl_3$  whole square plus  $dl_3, dl_1$  whole square.

### Cauchy Stress Components



The normal vector on area  $ABC = \vec{n} = n_1 \vec{e}_1 + n_2 \vec{e}_2 + n_3 \vec{e}_3$

$$= \frac{A_{PBC}}{A_{ABC}} \vec{e}_1 + \frac{A_{PAB}}{A_{ABC}} \vec{e}_2 + \frac{A_{PAC}}{A_{ABC}} \vec{e}_3 = \frac{dl_2 \cdot dl_3}{2 A_{ABC}} \vec{e}_1 + \frac{dl_3 \cdot dl_1}{2 A_{ABC}} \vec{e}_2 + \frac{dl_1 \cdot dl_2}{2 A_{ABC}} \vec{e}_3$$

$\therefore$  Area vector on surface  $ABC = A_{ABC} \cdot \vec{n}$

$$= \frac{1}{2} (dl_2 dl_3 \vec{e}_1 + dl_3 dl_1 \vec{e}_2 + dl_1 dl_2 \vec{e}_3)$$

$$\frac{dl'_1}{dl_1} = \frac{dl'_2}{dl_2} = \frac{dl'_3}{dl_3} \Rightarrow dl'_1 = x \frac{dl_1}{dl_3}, \quad dl'_2 = x \frac{dl_2}{dl_3}$$

Volume of the tetrahedron  $= V_{PABC} = \int_0^{dl_3} \frac{1}{2} dl'_1 dl'_2 dx = \frac{dl_1 dl_2}{2 dl_3^2} \int_0^{dl_3} x^2 dx$

$$= \frac{1}{6} dl_1 dl_2 dl_3$$

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Moving further, Considering this  $A_{ABC}$  with a normal vector  $\vec{n}$ , where  $n_1, n_2, n_3$  are the direction cosines for this normal vector. These direction cosines can be written as the ratio of the projected area and the inclined area. So,  $\vec{n}$  is the unit normal on this inclined  $A_{ABC}$ . If you project this inclined area on  $x_1$  plane,

that would become  $PBC$ . This particular plane, the back face, this is basically equals to the projection of  $ABC$  on the  $x_1$  plane. Thus, area of projected  $PBC$  divided by inclined area  $ABC$  gives us the first direction cosine  $n_1$ . Similarly, taking the projection of inclined area on  $x_2$  plane and  $x_3$  plane. We can write  $n_2$  as  $A_{PAB}$  by  $A_{ABC}$  and  $n_3$  as  $A_{PAC}$  divided by inclined  $A_{ABC}$ .

Now writing the expressions of all 3 areas of 3  $x_1, x_2, x_3$  planes  $PBC, PAB$  and  $PAC$  which we had already derived as half of  $dl_2, dl_3$  and so on. Substituting those we can get the normal vector. And going back to the previous slide if you see  $\vec{AB}$  cross  $\vec{BC}$  was defined as minus of 2 times inclined area dot  $\vec{n}$  vector. Now, this dot  $\vec{n}$  vector we had obtained thus the surface area vector for  $ABC$  can be written as half of  $dl_2 dl_3 \vec{e}_1$  plus  $dl_3 dl_1 \vec{e}_2$  plus  $dl_1 dl_2 \vec{e}_3$ . This gives us the surface area vector for this inclined surface  $ABC$ .

Now for finding the volume, let us take up another small slice which is shown here at a distance  $x$  from this tip  $B$ . then at a distance  $dx$  further from this  $x$ , we are taking another slice. And we will try to express the volume of this small region of the tetrahedron, which is sliced by drawing two planes at a distance  $x$  and  $x$  plus  $dx$  from the tip point  $B$ . So, considering the lengths at that particular cutting plane to be  $dl_1 dl_2$  and  $dl_3$  along the  $x_1, x_2, x_3$  axis,  $x_1, x_2, x_3$  directions and using the property of the symmetry of geometric symmetry we can write  $dl'_1$  prime by  $dl_1$  equals to  $dl'_2$  by  $dl_2, dl'_3$  prime by  $dl_3$  and using that  $x$  quantity  $dl'_1$  is equals to  $x$  times  $dl_1$  by  $dl_3$  because  $x$  is nothing but  $dl'_3$ .  $dl'_2$

can be written as  $x$  times  $dl_2$  by  $dl_3$  and volume of the small tetrahedron this small region elemental part of the tetrahedron can be written as half of  $dl_2' dl_3'$ . If you integrate it over the total length from  $B$  to  $P$ . So, volume of that small tetrahedron with thickness  $dx$  was half of  $dl_1' dl_2' dx$  integrating it for the total length of  $x$  varying from 0 to  $dl_3$  from point  $B$  to  $P$  we can get the total volume as  $x$  integral to be  $1$  by  $6 dl_1 dl_2 dl_3$ .

So, with this we got the volume of the tetrahedron we got the surface area vector of this inclined face  $ABC$ .

**Cauchy Stress Components**

The equations of motion along all three directions are obtained as:

$$t_1 A_{ABC} - \sigma_{11} A_{PBC} - \sigma_{21} A_{PAB} - \sigma_{31} A_{PAC} + b_1 \rho V_{FABC} = \rho a_1 V_{FABC}$$

$$t_2 A_{ABC} - \sigma_{12} A_{PBC} - \sigma_{22} A_{PAB} - \sigma_{32} A_{PAC} + b_2 \rho V_{FABC} = \rho a_2 V_{FABC}$$

$$t_3 A_{ABC} - \sigma_{13} A_{PBC} - \sigma_{23} A_{PAB} - \sigma_{33} A_{PAC} + b_3 \rho V_{FABC} = \rho a_3 V_{FABC}$$

The above equations are valid for both static and dynamic equilibriums

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Now, moving forward to the force balance or writing the governing equation along 1, 2 and 3 directions. The equations of motion along all 3 directions are obtained. So, along the  $x_1$  direction using the Newton's law of motion, total mass into acceleration is net force acting along a specific direction. So, considering  $a_1$  to be acceleration along the 1 direction, multiplied with the total mass mass is density times total volume so  $\rho$  times volume of the tetrahedron is total mass multiplied with one acceleration component along  $x$  direction this is mass into acceleration along  $x_1$  direction now left hand side is the net force along the  $x_1$  direction now If you consider the tetrahedron, there are four planes,  $ABC$ ,  $PBC$ ,  $PAC$  and  $PAB$ . On all four faces, we are having four traction vectors acting  $\tilde{t}^{(n)}$ ,  $t_1$ ,  $t_2$ ,  $t_3$  and each of them are having one component along  $x_1$ , another component along  $x_2$ , another component along  $x_3$ . So, taking the respective  $x_1$  components as this is the force balance along  $x_1$  direction,

We are taking the  $x_1$  component acting along all this. So, let us say  $ABC$  inclined phase  $x_1$  component is  $t_1$  multiplied with the corresponding area because we are writing total force. this fraction or this  $\sigma_{11}$  these are having these are stresses with the unit of force per unit area. So, we must multiply that with the corresponding area component to get the net

force. So,  $t_1 A_{ABC}$  is the net force acting on inclined face  $ABC$  along  $x_1$  direction minus  $\sigma_{11}$  area of  $PBC$  is net force acting along the  $x_1$  direction of  $PBC$ .

Similarly,  $\sigma_{21}$  area of  $PAB$  is force on the  $PAB$  plane along  $x_1$  direction minus  $\sigma_{31} A_{PAC}$  is the force acting on  $PAC$  side  $PAC$  plane along the 1 direction. So, these 4 are the forces along 1 direction for 4 different sides of the tetrahedron, 4 different planes of the tetrahedron. Now, coming to this term, this is defined  $b_1$  is defined to be the body force per unit mass. So, if you multiply that with the total mass  $\rho$  times volume, we will get the net body force along the 1 or  $x_1$  direction.

So, this completes the net equation of motion. Mass into acceleration along  $x_1$  direction is net force along  $x_1$  direction which includes surface fractions on four surfaces of the tetrahedron plus the body force within the body within the tetrahedron along  $x_1$  direction. Similarly, if you write for  $x_2$  and  $x_3$  directions, these would be the equations of motion coming for all three directions. So, these are the three equations of motion for  $x_1$ ,  $x_2$  and  $x_3$  direction.

Now, these equations are valid for both static and dynamic equilibrium because we are considering the acceleration of the system as well. Acceleration components were taken to be non-zero. So, these equations are valid for dynamic cases as well.

**Cauchy Stress Components**

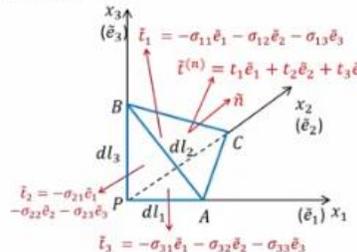
Dividing equations of motion by  $A_{ABC}$  and assuming static condition ( $a_i = 0$ ),

$$t_1 - \sigma_{11}n_1 - \sigma_{21}n_2 - \sigma_{31}n_3 + (b_1 - a_1) \frac{\rho V_{PABC}}{A_{ABC}} = 0$$

$$t_2 - \sigma_{12}n_1 - \sigma_{22}n_2 - \sigma_{32}n_3 + (b_2 - a_2) \frac{\rho V_{PABC}}{A_{ABC}} = 0$$

$$t_3 - \sigma_{13}n_1 - \sigma_{23}n_2 - \sigma_{33}n_3 + (b_3 - a_3) \frac{\rho V_{PABC}}{A_{ABC}} = 0$$

$$\frac{V_{PABC}}{A_{ABC}} = \frac{dl_1 dl_2 dl_3}{3\sqrt{(dl_1 dl_2)^2 + (dl_2 dl_3)^2 + (dl_3 dl_1)^2}}$$






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Now, moving forward. dividing all the terms of those three equations by the inclined area  $ABC$  and assuming the static condition.

So, previous expressions were valid for both static and dynamic. Now, moving to the static case, we are forcing acceleration components to be 0, forcing them to be 0 and dividing all the terms of three equations by inclined area  $ABC$ , we will be getting these three equations. Now, in the last term of all three equations, you can see  $\rho$  into volume by inclined area  $ABC$ . This term is coming.

So, from geometry, we had already derived the volume expression and inclined area  $ABC$  expression. So, this ratio, volume of the tetrahedron divided by inclined area  $ABC$  can be obtained like this. Now, as we are considering a small elemental tetrahedron with  $dl_1$ ,  $dl_2$ ,  $dl_3$  to be small, this particular quantity volume by area for the tetrahedron should also be a very small quantity. And with  $dl_1$ ,  $dl_2$ ,  $dl_3$  tending to 0 for small tetrahedron, this volume by area term, this entire term will go to 0. All the last terms in all three equations will go to 0.

**Cauchy Stress Components**

For infinitesimal tetrahedron,  $dl_1, dl_2, dl_3 \rightarrow 0 \Rightarrow \frac{V_{PABC}}{A_{ABC}} \rightarrow 0$

Thus, the force balance equations reduce to

$$t_1 = \sigma_{11}n_1 + \sigma_{21}n_2 + \sigma_{31}n_3$$

$$t_2 = \sigma_{12}n_1 + \sigma_{22}n_2 + \sigma_{32}n_3$$

$$t_3 = \sigma_{13}n_1 + \sigma_{23}n_2 + \sigma_{33}n_3$$

$$\Rightarrow \{t\} = [\sigma]^T \{n\}$$

Relation for obtaining traction vector acting on a plane at any point  $P$  with normal vector components  $\bar{n}$  in terms of Cauchy stress components.

$$\Rightarrow t_i = \sigma_{ij}^T n_j$$

$$\Rightarrow t_i = \sigma_{ji} n_j$$

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So, for infinitesimal small tetrahedron  $dl_1$ ,  $dl_2$ ,  $dl_3$  tending to 0, volume by area term goes to 0 and thus that static force balance reduces to this. The first equation along  $x_1$  becomes  $t_1$  equals to  $\sigma_{11}n_1$  plus  $\sigma_{21}n_2$  plus  $\sigma_{31}n_3$ . The expression along 2 direction  $x_2$  direction becomes  $t_2$  is  $\sigma_{12}n_1$  plus  $\sigma_{22}n_2$  plus  $\sigma_{32}n_3$ .  $t_3$  directional equation becomes  $t_3$  equals to  $\sigma_{13}n_1$  plus  $\sigma_{23}n_2$  plus  $\sigma_{33}n_3$ . Now, in total, using a single matrix notation, you can write the traction vector  $\tilde{t}$  to be  $\sigma$  transpose, transpose of the Cauchy stress tensor times  $n$  vector, where  $n$  is the unit normal vector on that particular plane. Now, as we had already taken  $dl_1$ ,  $dl_2$ ,  $dl_3$  tending to 0, this  $ABC$  plane is moving towards point  $P$  and we can consider that  $n$  is the unit normal vector of the plane passing through point  $P$  because of small tetrahedron assumption with  $dl_1$   $dl_2$   $dl_3$  tending to 0 volume by area ratio also went to 0. so this particular expression  $\{t\}$  equals to  $[\sigma]^T \{n\}$  this gives us the relation between the traction vector acting on a plane at a point  $P$  with a normal vector

direction  $n$  in terms of its cauchy stress components so you must remember that for defining the cauchy stress component we need a point and a plane both are required

plane is defined with respect to unit normal vector  $\tilde{n}$  or  $\tilde{n}$  vector and then only we can define the relation between cauchy stress components and surface traction vector as  $t$  equals to  $[\sigma]^T \{n\}$  In the indicial notation, we can write  $t_i$  to be  $\sigma_{ij}^T n_j$  or  $t_i$  to be  $\sigma_{ji} n_j$ . So, this relates the stress components, Cauchy stress components with the surface fraction component  $t_i$  equals to  $\sigma_{ji} n_j$ . So, in this lecture, we had discussed about the introduction to the surface traction, then Cauchy stress components and their relation with the help of the concept of Cauchy tetrahedron. Thank you.