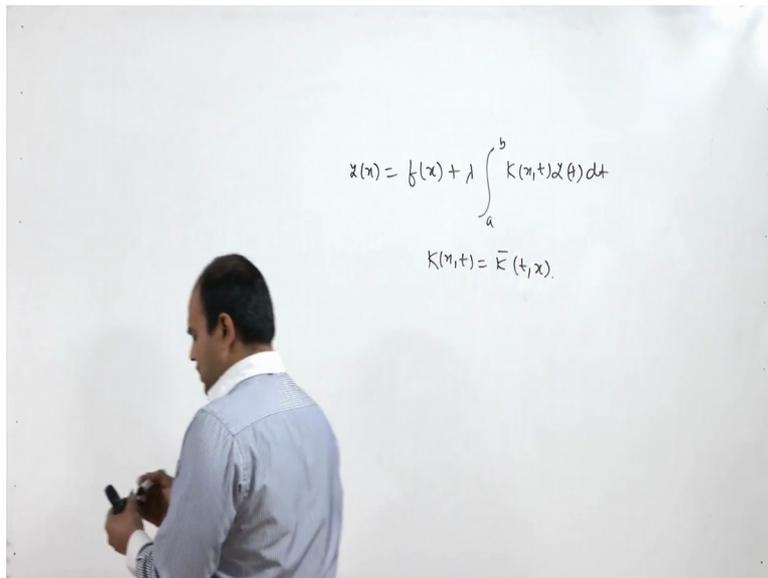


Integral Equations, Calculus of Variations and their Applications
Professor Dr. D. N. Pandey
Department of Mathematics
Indian Institute of Technology Roorkee
Lecture 13
Fredholm Integral Equations with Symmetric Kernels: Examples

Hello friends welcome to the lecture of Integral Equation and Calculus of Variations with its applications and examples, we were considering the case of symmetric kernel.

(Refer Slide Time: 0:38)



So if you remember we have we were doing this problem y of x equal to f of x plus λ a to b K of x, t y of t y of t d of t provided that this K of x, t is equal to \bar{K} t of x . So here we are assuming that we have a symmetric kernel and then we try to solve the Fredholm integral equation of second kind with this thing. So for that we have already discussed theory corresponding to eigenvalues and eigenfunction and how to solve this kind of equation, let us apply that theory to solve our problems.

(Refer Slide Time: 1:28)

Example 1

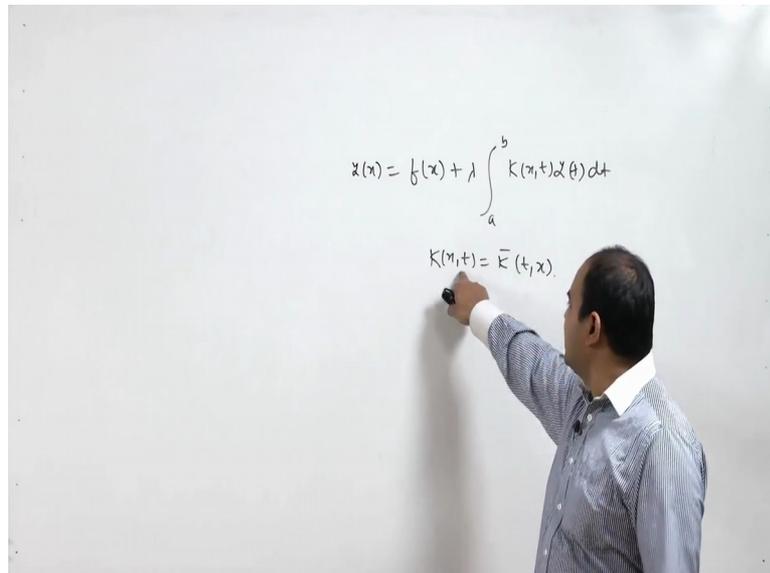
Solve the symmetric integral equation

$$y(x) = (x + 1)^2 + \int_{-1}^1 (xt + x^2 t^2) y(t) dt. \quad (20)$$

Here $f(x) = (x + 1)^2$ and $\lambda = 1$. Now we find the eigenvalues and corresponding eigenvectors to the homogeneous equation

$$\begin{aligned} y(x) &= \lambda \int_{-1}^1 (xt + x^2 t^2) y(t) dt \\ &= \lambda x \int_{-1}^1 t y(t) dt + \lambda x^2 \int_{-1}^1 t^2 y(t) dt \\ &= \lambda C_1 x + \lambda C_2 x^2, \end{aligned} \quad (21)$$

IT ROORKEE NPTEL ONLINE CERTIFICATION COURSE 16



So our first problem is this y of x equal to x plus 1 whole square plus a to b , here a is minus 1 and b is 1 xt plus x square t square y t dt . Now if you look at here your f of x is your x plus 1 whole square and λ is basically 1 here and if you look at K x , t is this so here K x , t is same as K t , x , here we are considering real so bar is not playing any role. So K x , t is equal to K t , x for this particular example.

If you remember this is of one more type that is separable type, so you can solve with the help of separable type also but we are solving this with the help of theory which we have developed earlier. So here we first we need to find out say eigenvalues and eigenfunction then to find out eigenvalues and eigenfunction let us consider the homogeneous version homogeneous version means this f of x is not there so we are considering this homogeneous

problem that is y of x equal to $\lambda - 1$ to 1 x t plus x square t square y t dt and we want to find out the values of λ for which we have a nontrivial solution and those that values those values we are saying that it is eigenvalues and corresponding solution eigenfunction.

So we are solving this with the help of method which we have developed for separable kernel, so we are writing this as $\lambda x - 1$ to 1 t y t dt plus λx square $- 1$ to 1 t square y t dt . Now this quantity is something which is a constant call it C_1 and this quantity $- 1$ to 1 t square y t is another constant which we call as C_2 , so y x is written as $\lambda C_1 x$ plus $\lambda C_2 x$ square, so solution we already know only thing is we need to find out this C_1 and C_2 . So for that we use the expression for C_1 and C_2 and use the value of y of x .

(Refer Slide Time: 3:24)

where

$$C_1 = \int_{-1}^1 ty(t)dt, \quad C_2 = \int_{-1}^1 t^2 y(t)dt. \quad (22)$$

Substituting the value of $y(t)$ from (21) in (22), we obtain

$$C_1 \left(1 - \frac{2\lambda}{3}\right) + 0 \cdot C_2 = 0, \quad (23)$$

$$0 \cdot C_1 + \left(1 - \frac{2\lambda}{5}\right) C_2 = 0. \quad (24)$$

IT ROORKEE NPTEL ONLINE CERTIFICATION COURSE 17

Example 1

Solve the symmetric integral equation

$$y(x) = (x+1)^2 + \int_{-1}^1 (xt + x^2 t^2) y(t) dt. \quad (20)$$

Here $f(x) = (x+1)^2$ and $\lambda = 1$. Now we find the eigenvalues and corresponding eigenvectors to the homogeneous equation

$$\begin{aligned} y(x) &= \lambda \int_{-1}^1 (xt + x^2 t^2) y(t) dt \\ &= \lambda x \int_{-1}^1 t y(t) dt + \lambda x^2 \int_{-1}^1 t^2 y(t) dt \\ &= \lambda C_1 x + \lambda C_2 x^2, \end{aligned} \quad (21)$$



Equations (23) and (24) have a nontrivial solution only if

$$D(\lambda) = \begin{vmatrix} 1 - \frac{2\lambda}{3} & 0 \\ 0 & 1 - \frac{2\lambda}{5} \end{vmatrix} = 0$$

which gives the eigenvalues $\lambda_1 = 3/2$ and $\lambda_2 = 5/2$. Substituting $\lambda = \lambda_1 = \frac{3}{2}$ in (23) and (24), we obtain

$$C_1 \cdot 0 + 0 \cdot C_2 = 0, \quad \text{and} \quad 0 \cdot C_1 + \left[1 - \left(\frac{2}{5} \times \frac{3}{2} \right) \right] C_2 = 0,$$

which gives $C_2 = 0$ and C_1 is arbitrary. Putting these values in (21), we obtain the eigenfunction

$$y_1(x) = \frac{3}{2} C_1 x.$$



So here we have C_1 is equal to this and C_2 is equal to this, so using $y(t)$ given by (21) you can get two equations (23) and (24), $C_1(1 - \frac{2\lambda}{3}) + 0 \cdot C_2 = 0$ and this thing. Now if you remember by solving this is nothing but algebraic equation ($D(\lambda) C_1$) this $D(\lambda)$ of C_1 and C_2 equal to this thing. So here your $D(\lambda)$ is the coefficient matrix of this algebraic equation.

So $D(\lambda)$ is coming out to be modulus of $1 - \frac{2\lambda}{3}$, 0 , 0 , $1 - \frac{2\lambda}{5}$ and if you look at this is a simply a triangular matrix and here you can get your eigenvalues very easily and you can say that $\lambda_1 = 3/2$ and $\lambda_2 = 5/2$ is the eigenvalues means for which your $D(\lambda)$ is having 0 determinant. So here eigenvalues we know then we need to find out the corresponding normalized orthonormalized eigenfunction.

So for that let us start with first that is lambda equal to lambda 1 that is 3 by 2 and try to solve this 23 and 24. So when you put lambda 1 equal to 3 by 2 here in 23 and 24 then you look at that the first equation 23 is simply redundant because this quantity is 0, so 0 times C 1 plus 0 times C 2 equal to 0, so this is simply will not give anything. So if you look at here by putting this you can get that C 2 equal to 0 and C 1 is completely arbitrary.

So when we have C 1 is completely arbitrary then you can write down our solution y 1 x here we can put the value of C 2, so C 2 is 0 here so y x is lambda C 1 x. So I can say that lambda is basically what lambda is 3 by 2 here, so I can write y 1 x is equal to 3 by 2 C 1 x. So corresponding to this eigenvalue lambda 1 equal to 3 by 2, we have a nontrivial solution that is 3 by 2 C 1 x.

Now with C 1 is some arbitrary constant so we can always take C 1 in a way such that this 3 by 2 C 1 is 1.

(Refer Slide Time: 5:38)

setting $\frac{3}{2}C_1 = 1$, we obtain $y_1(x) = x$ and the corresponding normalized eigenfunction $\psi_1(x)$ is given by

$$\psi_1(x) = \frac{y_1(x)}{\left[\int_{-1}^1 \{y_1(x)\}^2 dx \right]^{1/2}} = \frac{x}{\left[\int_{-1}^1 x^2 dx \right]^{1/2}} = \frac{\sqrt{6}x}{2}.$$

Now, substituting $\lambda = \lambda_2 = 5/2$ in (23) and (24), we obtain

$$\left[1 - \left(\frac{2}{3} \times \frac{5}{2} \right) \right] \cdot C_1 + 0 \cdot C_2 = 0, \quad \text{and} \quad 0 \cdot C_1 + 0 \cdot C_2 = 0,$$

which gives $C_1 = 0$ and C_2 is arbitrary. Putting these values in (21), we obtain the eigenfunction

$$y_2(x) = \frac{5}{2}C_2x^2$$




19

Equations (23) and (24) have a nontrivial solution only if

$$D(\lambda) = \begin{vmatrix} 1 - \frac{2\lambda}{3} & 0 \\ 0 & 1 - \frac{2\lambda}{5} \end{vmatrix} = 0$$

which gives the eigenvalues $\lambda_1 = 3/2$ and $\lambda_2 = 5/2$. Substituting $\lambda = \lambda_1 = \frac{3}{2}$ in (23) and (24), we obtain

$$C_1 \cdot 0 + 0 \cdot C_2 = 0, \quad \text{and} \quad 0 \cdot C_1 + \left[1 - \left(\frac{2}{5} \times \frac{3}{2} \right) \right] C_2 = 0,$$

which gives $C_2 = 0$ and C_1 is arbitrary. Putting these values in (21), we obtain the eigenfunction

$$y_1(x) = \frac{3}{2} C_1 x.$$

So we can say that your $y_1(x) = x$ is the eigenfunction corresponding to $\lambda_1 = 3/2$. Similarly we can find out (eigen) eigenfunction corresponding to $\lambda_2 = 5/2$. But we need normalized eigenfunction, so first let us normalize this $y_1(x)$, so normalize means you just divide by norm of this, so what is the norm here so norm here is $\int_{-1}^1 y_1(x)^2 dx$.

Now if you calculate $y_1(x) = x$ here, if you calculate this quantity is coming out to be $\frac{2}{\sqrt{6}}$, so you just calculate and divide by this, so $\Psi_1(x)$ is given by $\frac{\sqrt{6}}{2} x$. So this is very easy to solve, okay. So now λ_1 is known to you, Ψ_1 is known to you, now repeat the same procedure for λ_2 , so $\lambda_2 = 5/2$. So first equation gives you this and second is again the same way redundant.

So here you can solve and you can say that C_1 is coming out to be 0 and C_2 is arbitrary.

(Refer Slide Time: 6:48)

Example 1

Solve the symmetric integral equation

$$y(x) = (x+1)^2 + \int_{-1}^1 (xt + x^2 t^2) y(t) dt. \quad (20)$$

Here $f(x) = (x+1)^2$ and $\lambda = 1$. Now we find the eigenvalues and corresponding eigenvectors to the homogeneous equation

$$\begin{aligned} y(x) &= \lambda \int_{-1}^1 (xt + x^2 t^2) y(t) dt \\ &= \lambda x \int_{-1}^1 t y(t) dt + \lambda x^2 \int_{-1}^1 t^2 y(t) dt \\ &= \lambda C_1 x + \lambda C_2 x^2, \end{aligned} \quad (21)$$



setting $\frac{3}{2}C_1 = 1$, we obtain $y_1(x) = x$ and the corresponding normalized eigenfunction $\psi_1(x)$ is given by

$$\psi_1(x) = \frac{y_1(x)}{\left[\int_{-1}^1 \{y_1(x)\}^2 dx \right]^{1/2}} = \frac{x}{\left[\int_{-1}^1 x^2 dx \right]^{1/2}} = \frac{\sqrt{6}x}{2}.$$

Now, substituting $\lambda = \lambda_2 = 5/2$ in (23) and (24), we obtain

$$\left[1 - \left(\frac{2}{3} \times \frac{5}{2} \right) \right] C_1 + 0.C_2 = 0, \quad \text{and} \quad 0.C_1 + 0.C_2 = 0,$$

which gives $C_1 = 0$ and C_2 is arbitrary. Putting these values in (21), we obtain the eigenfunction

$$y_2(x) = \frac{5}{2} C_2 x^2$$



setting $\frac{5}{2}C_2 = 1$, we obtain $y_2(x) = x^2$ and the corresponding normalized eigenfunction $\psi_2(x)$ is given by

$$\psi_2(x) = \frac{y_2(x)}{\left[\int_{-1}^1 \{y_2(x)\}^2 dx \right]^{1/2}} = \frac{x^2}{\left[\int_{-1}^1 x^4 dx \right]^{1/2}} = \frac{\sqrt{10}x^2}{2}.$$

Now

$$f_1 = \int_{-1}^1 f(x) \psi_1(x) dx = \int_{-1}^1 (x+1)^2 \frac{\sqrt{6}x}{2} dx = \frac{2\sqrt{6}}{3} \quad (25)$$

$$\text{and } f_2 = \int_{-1}^1 f(x) \psi_2(x) dx = \int_{-1}^1 (x+1)^2 \frac{\sqrt{10}x^2}{2} dx = \frac{8\sqrt{10}}{15} \quad (26)$$



Again using the solution this $21 C_1$ is simply 0, so $y(x)$ is nothing but $\lambda C_2 x^2$. So here λ is $5/2$, so it is $y(x) = 5/2 C_2 x^2$. Now again in a same way C_2 is completely arbitrary, so we can always consider C_2 as 1 and so $y(x) = x^2$ is the eigenfunction corresponding to eigenvalue $5/2$. Now you normalize it so when you normalize it you divide by norm of this, so norm is $\int_{-1}^1 x^2 dx$, so $y(x) = x^2$ simply put it here and then you can find out the $\psi_2(x)$ under $\sqrt{10/2} x^2$.

(Refer Slide Time: 7:40)

Substituting this value of a_m in (16), we obtain the solution of the integral equation (15) in the form of an absolutely and uniformly convergent series

$$y(x) = f(x) + \lambda \sum_{m=1}^{\infty} \frac{f_m}{\lambda_m - \lambda} \psi_m(x)$$

$$\text{or } = f(x) + \lambda \sum_{m=1}^{\infty} \int \frac{\psi_m(x) \psi_m^*(t)}{\lambda_m - \lambda} f(t) dt$$

where the resolvent kernel is given by

$$\Gamma(x, t; \lambda) = \sum_{m=1}^{\infty} \frac{\psi_m(x) \psi_m^*(t)}{\lambda_m - \lambda}$$

Here we note that, the singular points of the resolvent kernel Γ corresponding to a symmetric L_2 -kernel are simple poles and every pole is an eigenvalue of the kernel.

IT ROORKEE NPTEL ONLINE CERTIFICATION COURSE 15

setting $\frac{5}{2} C_2 = 1$, we obtain $y_2(x) = x^2$ and the corresponding normalized eigenfunction $\psi_2(x)$ is given by

$$\psi_2(x) = \frac{y_2(x)}{\left[\int_{-1}^1 \{y_2(x)\}^2 dx \right]^{1/2}} = \frac{x^2}{\left[\int_{-1}^1 x^4 dx \right]^{1/2}} = \frac{\sqrt{10} x^2}{2}$$

Now

$$f_1 = \int_{-1}^1 f(x) \psi_1(x) dx = \int_{-1}^1 (x+1)^2 \frac{\sqrt{6} x}{2} dx = \frac{2\sqrt{6}}{3} \quad (25)$$

$$\text{and } f_2 = \int_{-1}^1 f(x) \psi_2(x) dx = \int_{-1}^1 (x+1)^2 \frac{\sqrt{10} x^2}{2} dx = \frac{8\sqrt{10}}{15} \quad (26)$$

IT ROORKEE NPTEL ONLINE CERTIFICATION COURSE 20

So once we have ψ_1 and ψ_2 then you can use our formula which says that our solution I am just going back to this and here we say that the solution is written as $y(x) = f(x) + \lambda \sum_{m=1}^{\infty} \frac{f_m}{\lambda_m - \lambda} \psi_m(x)$, here. Here m

is from 1 to infinity but here in our particular example we have only two eigenvalue so this infinite sum is reduced to only finite sum that is m equal to 1 and 2.

So what we want? We already have lambda 1 and lambda 2 and Psi 1 and Psi 2 we have, only thing we have to calculate is f 1 and f 2. So to calculate f 1 and f 2, f 1 and f 2 are fourier coefficient corresponding to f, so (f of) so f 1 is basically what? Minus 1 to 1 f of x Psi 1 x dx, so Psi 1 x is already known to you, you put it here and f x is already known to you so you can get the value of f 1 that is 2 root 6 by 3, this quite simple calculation you can just plug in and then you can find out f 1.

Similarly in a same way you can find out f of 2, f of 2 is minus 1 to 1 f of x Psi 2 x, so this is the fourier coefficient corresponding to this Psi 2. So again you can put it all values already you know f, you know Psi 2, you can get your f 2, f 2 is coming out to be 8 root 10 by 15.

(Refer Slide Time: 9:02)

Here $\lambda = 1$, $\lambda \neq \lambda_1 = \frac{3}{2}$ and $\lambda \neq \lambda_2 = \frac{5}{2}$. Therefore (20) will possess a unique solution given by

$$y(x) = f(x) + \lambda \sum_{m=1}^2 \frac{f_m}{\lambda_m - \lambda} \psi_m(x)$$

$$= (x+1)^2 + \frac{f_1 \psi_1(x)}{\lambda_1 - 1} + \frac{f_2 \psi_2(x)}{\lambda_2 - 1}$$

$$= 25/9x^2 + 6x + 1.$$

So now you know lambda 1, lambda 2, Psi 1 and Psi 2 and f 1 and f 2, you just write down your solution as y x equal to f x plus lambda m equal to 1 to 2 f m upon lambda m minus lambda Psi m x, so here lambda is 1 we have taken as 1, so f x is this, f 1 is this f 1 Psi 1 x upon lambda 1 minus 1 plus f 2 Psi 2 x upon lambda 2 minus 1, so putting all the values we have y x is equal to 25 by 9 x square plus 6 x plus 1.

(Refer Slide Time: 9:40)

Example 2

Solve the symmetric integral equation

$$y(x) = 1 + \lambda \int_0^{\pi} \cos(x+t)y(t)dt. \quad (27)$$

Here $f(x) = 1$ and $\lambda = \lambda$. Now we find the eigenvalues and corresponding eigenvectors to the homogeneous equation

$$y(x) = \lambda \int_0^{\pi} \cos(x+t)y(t)dt$$
$$= \lambda \int_0^{\pi} (\cos x \cos t - \sin x \sin t)y(t)dt \quad (28)$$
$$= \lambda \cos x \int_0^{\pi} \cos ty(t)dt - \lambda \sin x \int_0^{\pi} \sin ty(t)dt$$
$$= \lambda \cos x C_1 - \lambda \sin x C_2. \quad (29)$$

IT ROOKIEE NPTEL ONLINE CERTIFICATION COURSE 22

So let us consider one more example to have a better understanding I hope this is not very difficult. Let us look at this problem $y(x) = 1 + \lambda \int_0^{\pi} \cos(x+t)y(t)dt$. Now here again you look at $\cos(x+t)$ is a symmetric kernel, so it means that if you interchange it is first of all real then if you interchange the role of x and t then $\cos(x+t)$ is same as $\cos(t+x)$. So it is case of symmetric kernel and then we try to find out.

So here $f(x)$ is simply 1 and since λ is not given any constant value so we can say that λ is same as λ . So to find out the solution in terms of resolvent kernel or using our theory we need to calculate our eigenvalues and eigenfunction. So consider the homogeneous equation that is you take simply $f(x)$ as 0, so consider this $y(x) = \lambda \int_0^{\pi} \cos(x+t)y(t)dt$, now we want to find out say eigenvalues and eigenvectors.

So this is this you can solve using the method of separable kernel, so you can find out eigenvalues and eigenfunction with the help of method we have discussed for separable kernel. So here you write down $\cos(x+t)$ as $\cos x \cos t - \sin x \sin t$. Now again you simplify then it is $\int_0^{\pi} \cos ty(t)dt - \lambda \sin x \int_0^{\pi} \sin ty(t)dt$. Now this quantity is again constant call it C_1 , this quantity is your constant call it C_2 .

So it means that your $y(x)$ is already known to you, so this is $\lambda \cos x C_1 - \lambda \sin x C_2$. So only you need to find out this C_1 and C_2 .

(Refer Slide Time: 11:24)

where

$$C_1 = \int_0^\pi \cos ty(t)dt \quad \text{and} \quad C_2 = \int_0^\pi \sin ty(t)dt \quad (30)$$

Substituting the value of $y(t)$ from (29) in (30), we obtain

$$C_1(2 - \lambda\pi) + 0.C_2 = 0, \quad (31)$$
$$0.C_1 + (2 + \lambda\pi)C_2 = 0. \quad (32)$$


IT ROORKEE NPTEL ONLINE CERTIFICATION COURSE

Example 2

Solve the symmetric integral equation

$$y(x) = 1 + \lambda \int_0^\pi \cos(x+t)y(t)dt. \quad (27)$$

Here $f(x) = 1$ and $\lambda = \lambda$. Now we find the eigenvalues and corresponding eigenvectors to the homogeneous equation

$$y(x) = \lambda \int_0^\pi \cos(x+t)y(t)dt$$
$$= \lambda \int_0^\pi (\cos x \cos t - \sin x \sin t)y(t)dt \quad (28)$$
$$= \lambda \cos x \int_0^\pi \cos ty(t)dt - \lambda \sin x \int_0^\pi \sin ty(t)dt$$
$$= \lambda \cos x C_1 - \lambda \sin x C_2. \quad (29)$$


IT ROORKEE NPTEL ONLINE CERTIFICATION COURSE

Equations (31) and (32) have a nontrivial solution only if

$$D(\lambda) = \begin{vmatrix} 2 - \lambda\pi & 0 \\ 0 & 2 + \lambda\pi \end{vmatrix} = 0$$

which gives the eigenvalues $\lambda_1 = 2/\pi$ and $\lambda_2 = -2/\pi$. Substituting $\lambda = \lambda_1 = 2/\pi$ in (31) and (32), we obtain

$$0.C_1 + 0.C_2 = 0, \quad \text{and} \quad 0.C_1 + 4C_2 = 0$$

which gives $C_2 = 0$ and C_1 is arbitrary. Putting these values in (29), we obtain the eigenfunction

$$y_1(x) = \frac{2}{\pi} C_1 \cos x.$$


IT ROORKEE NPTEL ONLINE CERTIFICATION COURSE

So here using this we already know that C_1 is $\int_0^\pi \cos t y t dt$ and C_2 is $\int_0^\pi \sin t y t dt$. So C_1 and C_2 and you already know the solution y like this, so from this you can calculate $y t$ put it back to C_1 and C_2 you can have these two equations $C_1^2 - \lambda \pi + 0 \times C_2 = 0$ and this is equation number 32 and this is again if you look at this is again a triangular equation to solve and here your eigenvalue can be found out by say coefficient matrix determinant of coefficient matrix and determinant of coefficient matrix is nothing but $D(\lambda)$ and when you solve this you will get that we have two eigenvalues that is λ_1 is equal to 2 by π and λ_2 equal to -2 by π .

So eigenvalues is with us, now we try to find out the corresponding eigenfunction. So let us start with λ equal to λ_1 that is 2 by π , so when you put λ equal to λ_1 that is 2 by π then first equation is simply redundant and you can solve this and you can say that C_2 is coming out to be 0 and C_1 is arbitrary as we have considered in the previous case.

So here we say that C_2 is 0 and C_1 is arbitrary and which says that look at equation number 29. So $y(x)$ is already given so using $C_2 = 0$ and C_1 arbitrary $y_1(x)$ is equal to 2 by $\pi C_1 \cos$ of x . Again in a same way which we have already discussed you can take 2 by πC_1 as 1 , C_1 is arbitrary.

(Refer Slide Time: 13:08)

setting $\frac{2}{\pi} C_1 = 1$, we obtain $y_1(x) = \cos x$ and the corresponding normalized eigenfunction $\psi_1(x)$ is given by

$$\psi_1(x) = \frac{y_1(x)}{\left[\int_0^\pi \{y_1(x)\}^2 dx \right]^{1/2}} = \frac{\cos x}{\left[\int_0^\pi \cos^2 x dx \right]^{1/2}} = \left(\frac{2}{\pi} \right)^{1/2} \cos x.$$

Now, substituting $\lambda = \lambda_2 = -\frac{2}{\pi}$ in (31) and (32), we obtain

$$4C_1 + 0.C_2 = 0, \quad \text{and} \quad 0.C_1 + 0.C_2 = 0,$$

which gives $C_1 = 0$ and C_2 is arbitrary. Putting these values in (29), we obtain the eigenfunction

$$y_2(x) = -\left(\frac{-2}{\pi} \right) C_2 \sin x.$$

IFT ROORKEE NPTEL ONLINE CERTIFICATION COURSE 25

Example 2

Solve the symmetric integral equation

$$y(x) = 1 + \lambda \int_0^{\pi} \cos(x+t)y(t)dt. \quad (27)$$

Here $f(x) = 1$ and $\lambda = \lambda$. Now we find the eigenvalues and corresponding eigenvectors to the homogeneous equation

$$y(x) = \lambda \int_0^{\pi} \cos(x+t)y(t)dt$$

$$= \lambda \int_0^{\pi} (\cos x \cos t - \sin x \sin t)y(t)dt \quad (28)$$

$$= \lambda \cos x \int_0^{\pi} \cos ty(t)dt - \lambda \sin x \int_0^{\pi} \sin ty(t)dt$$

$$= \lambda \cos x C_1 - \lambda \sin x C_2. \quad (29)$$

IT ROORKEE NPTEL ONLINE CERTIFICATION COURSE 22

So $y = 1 + \lambda \cos x$ is your \cos of x is the corresponding eigenfunction, so just normalize it. So how to normalize it you divide it by norm norm here is a to b , a is what 0 and b is π , so I think it is already cleared here a is 0 and b is π , so $\int_0^{\pi} 1^2 dx = x \Big|_0^{\pi} = \pi$ so this is a norm of 1 so divide by this and this quantity is simply $\frac{1}{\sqrt{\pi}}$, so this is what $\frac{1}{\sqrt{\pi}}$ by \cos of x .

So this is the eigenvalue normalized eigenfunction corresponding to $\lambda = \frac{1}{\sqrt{\pi}}$. Now let us do the same procedure for $\lambda = -\frac{2}{\sqrt{\pi}}$ and you can say that in a same way you can get $C_1 = 0$ and C_2 is arbitrary and you can write down the corresponding nontrivial solution as $y = \sin x$ and it is minus of $\frac{2}{\sqrt{\pi}}$ $C_2 \sin x$, again you can take this constant as 1 and $y = \sin x$ is nothing but \sin of x is the corresponding eigenfunction.

(Refer Slide Time: 14:18)

setting $\frac{2}{\pi}C_2 = 1$, we obtain $y_2(x) = \sin x$ and the corresponding normalized eigenfunction $\psi_2(x)$ is given by

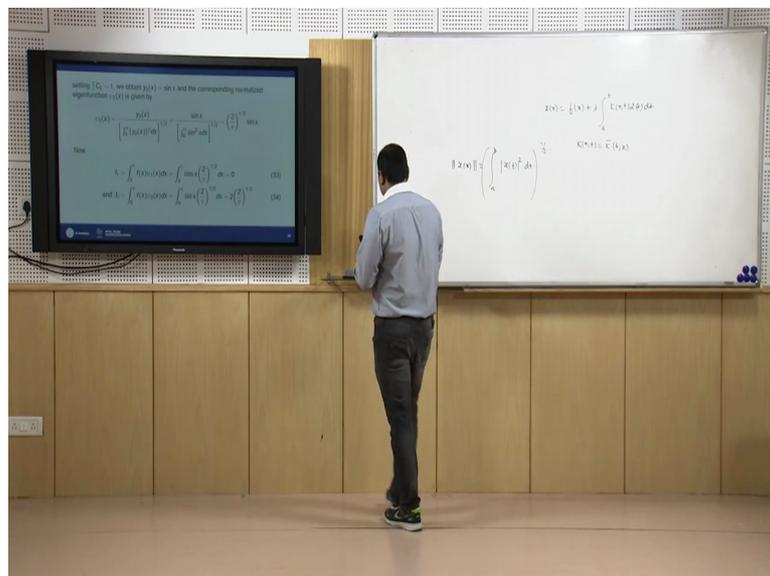
$$\psi_2(x) = \frac{y_2(x)}{\left[\int_0^\pi \{y_2(x)\}^2 dx\right]^{1/2}} = \frac{\sin x}{\left[\int_0^\pi \sin^2 x dx\right]^{1/2}} = \left(\frac{2}{\pi}\right)^{1/2} \sin x.$$

Now

$$f_1 = \int_0^\pi f(x)\psi_1(x)dx = \int_0^\pi \cos x \left(\frac{2}{\pi}\right)^{1/2} dx = 0 \quad (33)$$

$$\text{and } f_2 = \int_0^\pi f(x)\psi_2(x)dx = \int_0^\pi \sin x \left(\frac{2}{\pi}\right)^{1/2} dx = 2\left(\frac{2}{\pi}\right)^{1/2}. \quad (34)$$

IT ROORKEE NPTEL ONLINE CERTIFICATION COURSE 26



You can again normalize in a same way so $\psi_2(x)$ is $y_2(x)$ divided by norm of $y_2(x)$ which is nothing but this quantity. So please I hope we already know norm of y of x is equal to a to b of x square d of x 1 by 2, okay by the way this is this is a dummy variable so you can use t or anything, I am using t because I am using x here, okay. So this is norm of y of x so that is why we are dividing by norm of this. So here we have $\psi_2(x)$ is given to you.

Now in a similar way we can find out f_1 and f_2 which is fourier coefficient of f corresponding to these eigenfunction ψ_1 and ψ_2 , so f_1 is 0 to π a to b , so a to b means 0 to π $f(x)\psi_1(x)dx$ and you can say that it is coming out to be 0 and f_2 is 0 to π $f(x)\psi_2(x)dx$ which you can calculate it is coming out to be this quantity 2 upon 2 by π square root of this 2 into square root of 2 by π . So f_1, f_2 is known to us.

Now you simply put it back to your solution and now problem is that this lambda is not given any constant value. So your solution will also depend on the value of lambda because lambda we can take any constant, so right now we have a choice that lambda is none of the eigenvalues, one of the eigenvalues you can take this thing.

(Refer Slide Time: 16:06)

Now three cases arise:

Case 1: If $\lambda \neq \lambda_1$ and $\lambda \neq \lambda_2$. Then (27) will possess unique solution given by

$$\begin{aligned}
 y(x) &= f(x) + \lambda \sum_{m=1}^2 \frac{f_m}{\lambda_m - \lambda} \psi_m(x) \\
 &= 1 + \frac{\lambda}{\lambda_1 - \lambda} f_1 \psi_1(x) + \frac{\lambda}{\lambda_2 - \lambda} f_2 \psi_2(x) \\
 &= 1 - \frac{4\lambda \sin x}{2 + \lambda\pi}.
 \end{aligned}$$

Case 2: If $\lambda = \lambda_2 = -2/\pi$. Since $f_2 \neq 0$, so (27) possess no solution.





27

So let us consider the case where lambda is say taking no values taken as lambda 1 and lambda 2. So lambda is neither lambda 1 or lambda 2. So in this case you have y x equal to f of x plus lambda m equal to 1 to 2 f m lambda m minus lambda Psi m x. So writing f 1 and f 2 you can get the solution like this. Now if you look at here look at this 2 plus lambda pi, so if so in case when lambda is neither of these eigenvalues you can write y of x as this.

Now in a second case when lambda is equal to lambda 2, so if you look at lambda equal to lambda 2 here then this is something which creates problem. Now here lambda is lambda 2 means denominator is 0 and if you look at numerator lambda is some fixed quantity, now look at this Psi 2 is nonzero so only thing depend on this f 2 norm. Now if f 2 is again 0, if f 2 is coming out to be 0 then it is kind of 0 by 0 form, right.

And in this case this can take any value. So right now we say that f 2 is nonzero which we have calculated by equation number 34 that f 2 is nonzero. So in this case it is (0) nonzero divided by 0, so this cannot happen. So it means that in this case when lambda is equal to lambda 2 that is minus 2 by pi we do not have any solution. So it means that no solution exists if lambda has value minus 2 by pi.

(Refer Slide Time: 17:36)

Case 3: If $\lambda = \lambda_1 = 2/\pi$. Since $f_1 = 0$, there exists infinitely many solutions given by

$$\begin{aligned} y(x) &= f(x) + A\psi_1(x) + \frac{\lambda}{\lambda_m - \lambda} f_m \psi_m(x) \\ &= 1 + A \left(\frac{2}{\pi}\right)^{1/2} \cos x + \frac{(2/\pi)}{-(2/\pi) - (2/\pi)} \times 2 \left(\frac{2}{\pi}\right) \sin x \\ &= 1 + C \cos x - \frac{2}{\pi} \sin x. \end{aligned}$$

where $C = A \left(\frac{2}{\pi}\right)^{1/2}$ is an arbitrary constant.



setting $\frac{2}{\pi} C_2 = 1$, we obtain $y_2(x) = \sin x$ and the corresponding normalized eigenfunction $\psi_2(x)$ is given by

$$\psi_2(x) = \frac{y_2(x)}{\left[\int_0^\pi \{y_2(x)\}^2 dx\right]^{1/2}} = \frac{\sin x}{\left[\int_0^\pi \sin^2 x dx\right]^{1/2}} = \left(\frac{2}{\pi}\right)^{1/2} \sin x.$$

Now

$$f_1 = \int_0^\pi f(x) \psi_1(x) dx = \int_0^\pi \cos x \left(\frac{2}{\pi}\right)^{1/2} dx = 0 \quad (33)$$

$$\text{and } f_2 = \int_0^\pi f(x) \psi_2(x) dx = \int_0^\pi \sin x \left(\frac{2}{\pi}\right)^{1/2} dx = 2 \left(\frac{2}{\pi}\right)^{1/2}. \quad (34)$$



Now three cases arise:

Case 1: If $\lambda \neq \lambda_1$ and $\lambda \neq \lambda_2$. Then (27) will possess unique solution given by

$$\begin{aligned} y(x) &= f(x) + \lambda \sum_{m=1}^2 \frac{f_m}{\lambda_m - \lambda} \psi_m(x) \\ &= 1 + \frac{\lambda}{\lambda_1 - \lambda} f_1 \psi_1(x) + \frac{\lambda}{\lambda_2 - \lambda} f_2 \psi_2(x) \\ &= 1 - \frac{4\lambda \sin x}{2 + \lambda\pi}. \end{aligned}$$

Case 2: If $\lambda = \lambda_2 = -2/\pi$. Since $f_2 \neq 0$, so (27) possess no solution.



Now let us take λ equal to λ_1 that is 2 by π , now in this case f_1 is coming out to be 0 that is calculated here 33 , so it means that here your f_1 is 0 denominator is also 0 so it is kind of 0 by 0 form. So this I can take any value A so if this limit exist we call it any value 0 by 0 is you can take any value A you can take any value so here we have infinitely many solution so I can write it $A \Psi_1 x$ plus λ upon λ_m minus λ $f_m \Psi_m x$. So here λ_m is your 2 , okay.

So here I can write A is any arbitrary constant and this λ_m is λ_2 so you can put it the value here and you can say that it is 1 plus $C \cos x$ minus 2 by π \sin of x where C is this and since A is arbitrary, C is also an arbitrary thing. So in this particular case if λ is one of the eigenvalue and if the corresponding fourier coefficient is nonzero then we have no solution as we have discussed in case 2.

But if λ is one of the eigenvalue and the corresponding fourier coefficient f_1 say f_5 whatever is 0 , so it means that when λ takes eigenvalues and the f nonhomogeneous term is orthogonal to the corresponding eigenfunctions in that case we have infinitely many solution and this we can see in a this we have already discussed in the case when we have separable kernel. So in case of separable kernel this case we have already discussed for the example when we have taken kernel $K(x, t)$ as 1 minus $3xt$.

So here we are considering the same case that λ equal to λ_1 , so λ is taking the first eigenvalue and f_1 is 0 means that f is orthogonal to the first eigenfunction corresponding to the eigenvalue here. So if it is the case then we have infinitely many solution and if that λ is λ_2 and f is not orthogonal to the second eigenfunction then we do not have any solution. So we have these many cases.

(Refer Slide Time: 20:16)

Fredholm Integral Equation of First Kind

Consider the Fredholm integral equation of first kind

$$f(x) = \int K(x, t)g(t)dt, \quad (35)$$

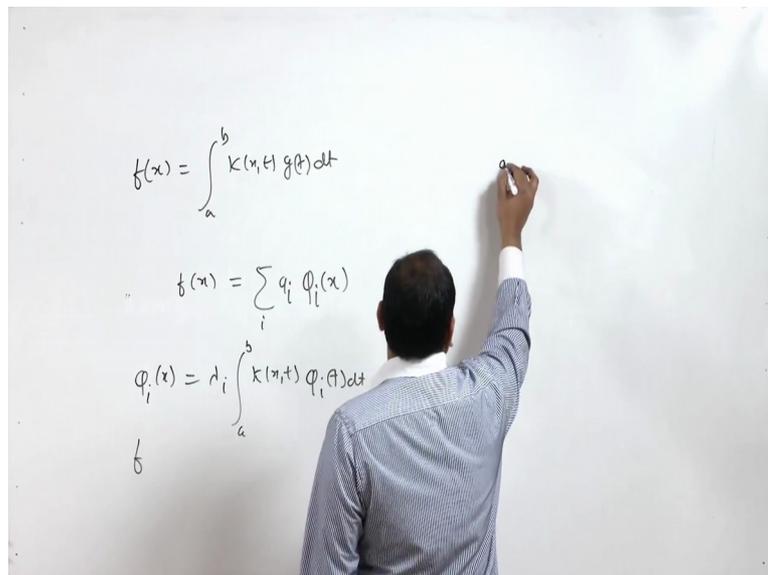
where $K(x, t)$ is a symmetric L_2 -kernel. Suppose $\lambda_1, \lambda_2, \dots, \lambda_k, \dots$ be the eigenvalues and $\psi_1, \psi_2, \dots, \psi_k, \dots$ be the corresponding eigenfunctions. Then from (14), we have

$$f_k = g_k / \lambda_k$$
$$g_k = \lambda_k f_k.$$

IT ROORKEE NPTEL ONLINE CERTIFICATION COURSE 29

Now with the help of this now let us consider this fredholm integral equation of the first kind. Here again we use Hilbert Schmidt theorem which says that f we have f x equal to K x, t is g t dt then where K and g both are L 2 functions. Then f x can be written as your linear combination of eigenfunctions corresponding to K x, t so that we already so Hilbert Schmidt theorem we are assuming.

(Refer Slide Time: 20:58)



Fredholm Integral Equation of First Kind

Consider the Fredholm integral equation of first kind

$$f(x) = \int K(x, t)g(t)dt, \quad (35)$$

where $K(x, t)$ is a symmetric L_2 -kernel. Suppose $\lambda_1, \lambda_2, \dots, \lambda_k, \dots$ be the eigenvalues and $\psi_1, \psi_2, \dots, \psi_k, \dots$ be the corresponding eigenfunctions. Then from (14), we have

$$\begin{aligned} f_k &= g_k / \lambda_k \\ g_k &= \lambda_k f_k. \end{aligned}$$

So here let me write it here so here we are using this is the Hilbert Schmidt theorem that f of x is equal to $\int_a^b K(x, t)g(t)dt$ and if K and g both are L_2 function then f can be written as summation of $a_i \psi_i(x)$ and summation over i and where $\psi_i(x)$ is normalized eigenfunction corresponding to this kernel K . So it means that we are solving this y of x is equal to $\int_a^b K(x, t)g(t)dt$ and $y(x)$, so here $\psi_i(x)$ is the solution of this, so let me write it here $\psi_i(x) = \lambda_i^{-1} f_i(x)$ okay.

And this Fourier coefficient here f_i is related to Fourier coefficient g_i like this, so here you can say that your f_k is written as g_k divided by λ_k , this we discussed earlier. So here I can write g_k as $\lambda_k f_k$. Now if you look at unknown function we are assuming in the first kind we are assuming that this g is the unknown function which we are discussing. So how to construct this unknown function from the Fourier coefficient of this unknown function g .

(Refer Slide Time: 22:24)

From Riesz-Fischer theorem, there are only two possibilities: either the infinite series

$$\sum_{k=1}^{\infty} f_k^2 \lambda_k^2 \quad (36)$$

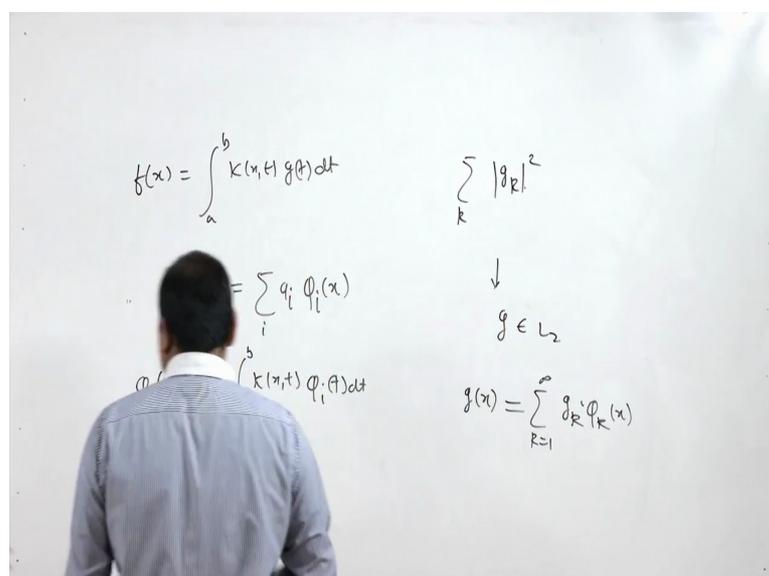
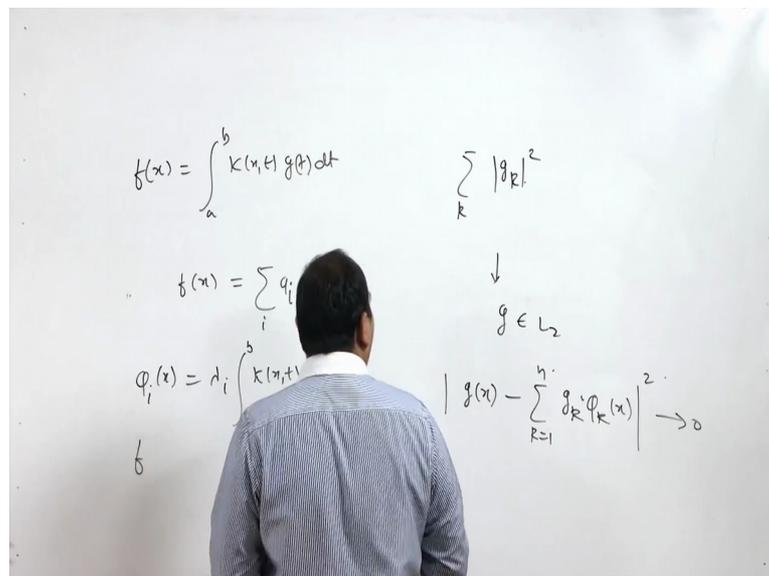
diverges and (35) has no solution, or the series (36) converges and there is a unique L_2 -function $g(x)$ which is the solution of (35) and given by

$$g(x) = \lim_{n \rightarrow \infty} \sum_{k=1}^n \lambda_k f_k v_k(x). \quad (37)$$



 NPTEL ONLINE CERTIFICATION COURSE

30



So here we have to recall the Riesz-Fischer theorem, which says that if we have say this g_k , right. So if this series is absolutely and uniformly convergent then we have a function g of course in L^2 such that this g has this representation $g(x) = \sum_{k=1}^{\infty} a_k \phi_k(x)$ and $\sum_{k=1}^n a_k^2 \rightarrow 0$, okay or you can say that that g has this representation, okay.

So so you can write that if this infinite series if it diverges then we do not have any solution, but if it converges then Riesz-Fischer theorem says that your solution $g(x)$ can be written as this limit of this infinite series, okay and this limit this Riesz-Fischer theorem says so. So it means that everything depend on this infinite series.

(Refer Slide Time: 24:12)

Example 3

Solve the symmetric Fredholm integral equation of first kind

$$f(x) = \int_0^1 K(x, t)g(t)dt, \quad (38)$$

where

$$K(x, t) = \begin{cases} x(1-t), & x < t; \\ (1-x)t, & x > t. \end{cases}$$

As we know that the boundary value problem

$$\frac{d^2y}{dx^2} + \lambda y = 0, \quad y(0) = y(1) = 0$$

is equivalent to the homogeneous equation

$$g(x) = \lambda \int_0^1 K(x, t)g(t)dt, \quad (39)$$

IT ROORKEE NPTEL ONLINE CERTIFICATION COURSE 31

Fredholm Integral Equation of First Kind

Consider the Fredholm integral equation of first kind

$$f(x) = \int K(x, t)g(t)dt, \quad (35)$$

where $K(x, t)$ is a symmetric L_2 -kernel. Suppose $\lambda_1, \lambda_2, \dots, \lambda_k, \dots$ be the eigenvalues and $\psi_1, \psi_2, \dots, \psi_k, \dots$ be the corresponding eigenfunctions. Then from (14), we have

$$\begin{aligned} \hat{f}_k &= g_k / \lambda_k \\ g_k &= \lambda_k \hat{f}_k. \end{aligned}$$

IT ROORKEE NPTEL ONLINE CERTIFICATION COURSE 29

So let us take one simple example here, so let us take this example here f of x equal to 0 to 1 $K x, t g t dt$ where $K x, t$ is given by this, so here it is the fredholm integral equation of the first kind. Then we need to find out that this is falling in this category I need to find out $f k$ and $g k$ all these thing. So eigenvalues we need to find out and eigenfunction we need to find out.

So for that corresponding to this we need to find out the eigenvalue and eigenproblem. So here for that we need to find out we need to solve this where $K x, t$ is given by this this thing.

(Refer Slide Time: 25:00)

$$f(x) = \int_a^b K(x,t) g(t) dt$$

$$f(x) = \sum_i q_i \phi_i(x)$$

$$\phi_i(x) = \lambda_i \int_a^b K(x,t) \phi_i(t) dt$$

$$\phi(x) = \lambda \int_a^b K(x,t) \phi(t) dt$$

$$\sum_k |q_k|^2$$

$$\downarrow$$

$$g \in L_2$$

$$g(x) = \sum_{k=1}^{\infty} g_k \phi_k(x)$$

So to find out this we can let me write it here ϕ of x equal to λ here it is 0 to 1 $K x, t$ and your $\phi t dt$, okay. So from this you can easily say that by conversion of this integral equation into corresponding boundary value problem we can say that the corresponding boundary value problem is this $d^2 y$ by dx square plus λy equal to 0 with the boundary condition that y of 0 equal to y of 1 equal to 0, so that is clear from this equation that y of 0 when you put it here so y of 0 equal to y 1 equal to 0.

(Refer Slide Time: 25:42)

Example 3

Solve the symmetric Fredholm integral equation of first kind

$$f(x) = \int_0^1 K(x, t)g(t)dt, \quad (38)$$

where

$$K(x, t) = \begin{cases} x(1-t), & x < t; \\ (1-x)t, & x > t. \end{cases}$$

As we know that the boundary value problem

$$\frac{d^2y}{dx^2} + \lambda y = 0, \quad y(0) = y(1) = 0$$

is equivalent to the homogeneous equation

$$g(x) = \lambda \int_0^1 K(x, t)g(t)dt, \quad (39)$$

IT ROORKEE NPTEL ONLINE CERTIFICATION COURSE 31

The eigenvalues of the system are given by $\lambda_1 = \pi^2, \lambda_2 = (2\pi)^2, \dots$ and the corresponding eigenfunctions are $\sqrt{2} \sin \pi x, \sqrt{2} \sin 2\pi x, \dots$. Therefore

$$f_k = \sqrt{2} \int_0^1 (\sin k\pi t)f(t)dt$$

and the integral equation has a solution of class L_2 if and only if the infinite series

$$\sum_{k=1}^{\infty} f_k^2 \lambda_k^2 = \pi^4 \sum_{k=1}^{\infty} (k^4 f_k^2)$$

converges.

IT ROORKEE NPTEL ONLINE CERTIFICATION COURSE 32

So to solve this we have to solve this particular problem, so I think this is quite simple problem and you can say that in this case your lambda i are simply what pi square and basically it is n pi square n pi whole square and the corresponding eigenfunctions are root 2 sin pi x root 2 sin 2 pi x and you can say that root 2 sin n pi x are your eigenfunction. So it means that corresponding to this kernel look at the eigenvalue problem that is g of x equal lambda 0 to 1 K x, t g t dt.

And you look at that this is equal and to the boundary value problem here, so find out the solution here. So whatever be the solutions here we have this solution here and the values for which we are getting a nontrivial solution here the same will be nontrivial solution here. So it

means that solving this problem is equivalent to solving this problem. So here we are finding the lambda for which we have a nontrivial solution.

(Refer Slide Time: 26:44)

The eigenvalues of the system are given by $\lambda_1 = \pi^2$, $\lambda_2 = (2\pi)^2, \dots$ and the corresponding eigenfunctions are $\sqrt{2} \sin \pi x$, $\sqrt{2} \sin 2\pi x, \dots$. Therefore

$$f_k = \sqrt{2} \int_0^1 (\sin k\pi t) f(t) dt$$

and the integral equation has a solution of class L_2 if and only if the infinite series

$$\sum_{k=1}^{\infty} f_k^2 \lambda_k^2 = \pi^4 \sum_{k=1}^{\infty} (k^4 f_k^2)$$

converges.

IT ROORKEE NPTEL ONLINE CERTIFICATION COURSE 32

From Riesz-Fischer theorem, there are only two possibilities: either the infinite series

$$\sum_{k=1}^{\infty} f_k^2 \lambda_k^2 \tag{36}$$

diverges and (35) has no solution, or the series (36) converges and there is a unique L_2 -function $g(x)$ which is the solution of (35) and given by

$$g(x) = \lim_{n \rightarrow \infty} \sum_{k=1}^n \lambda_k f_k \psi_k(x). \tag{37}$$

IT ROORKEE NPTEL ONLINE CERTIFICATION COURSE 30

So lambda is coming out to be n pi whole square and the corresponding eigenfunctions are this. So now you can calculate your fourier coefficient f k like this root 2 0 to 1 sin k pi t f t dt. Now once we have f k you can calculate f k and lambda k square. Now this is this quantity pi to power 4 k equal to 1 to infinity k 4 f k square. Now if this series converges then you can write down your g as this limit, so once this lambda f k this series is convergent then you can write yours g x as limit of this series.

So it means that we need to find out f_k such that this series is converges one this series is converges you can find out your solution g , is it okay. So we will discuss some more problem in our exercise so thank you for listening us we will meet in next lecture, thank you.