

# Lecture 04: Behavior of Regular Sturm–Liouville Systems

## Part II

### 1 The Sturm-Liouville Operator

Let  $I = (a, b)$  be a finite interval. We define the Sturm–Liouville operator

$$L[y] := (p(x)y'(x))' + q(x)y(x),$$

where the coefficient functions satisfy

$$p, q, s \in C^\infty([a, b]), \quad p(x) > 0, \quad s(x) > 0 \text{ on } [a, b].$$

For simplicity, we take the domain of  $L$  as

$$D(L) = C^\infty([a, b]),$$

though in general  $C^2([a, b])$  suffices.

The associated eigenvalue problem is

$$L[y] + \lambda s(x)y = 0, \tag{1}$$

together with regular Sturm–Liouville boundary conditions

$$\begin{aligned} \alpha_1 y(a) + \alpha_2 y'(a) &= 0, \\ \beta_1 y(b) + \beta_2 y'(b) &= 0, \end{aligned}$$

where  $\alpha_1^2 + \alpha_2^2 \neq 0$  and  $\beta_1^2 + \beta_2^2 \neq 0$ .

### 2 Orthogonality of Eigenfunctions

If  $(\lambda_k, \phi_k)$  and  $(\lambda_\ell, \phi_\ell)$  are two distinct eigenpairs with

$$\lambda_k \neq \lambda_\ell,$$

then the corresponding eigenfunctions are orthogonal with respect to the weight  $s(x)$ :

$$\int_a^b s(x)\phi_k(x)\phi_\ell(x) dx = 0.$$

### 3 Lagrange Identity

**Theorem 1** (Lagrange Identity). *Let  $y, z \in D(L)$ . Then*

$$yL[z] - zL[y] = \frac{d}{dx} \left( p(x)(yz' - zy') \right).$$

*Proof.* A direct computation shows

$$y(pz')' - z(py')' = \frac{d}{dx} (p(yz' - zy')),$$

and the  $qyz$  terms cancel. □

### 4 Self-Adjointness of the Sturm-Liouville Operator

We define the  $L^2$  inner product

$$\langle f, g \rangle := \int_a^b f(x) \overline{g(x)} dx.$$

**Theorem 2.** *The Sturm-Liouville operator  $L$  is self-adjoint, i.e.*

$$\langle L[y], z \rangle = \langle y, L[z] \rangle \quad \text{for all } y, z \in D(L).$$

*Proof.* Using the Lagrange identity,

$$\langle L[y], z \rangle - \langle y, L[z] \rangle = \left[ p(x)(\bar{z}y' - y\bar{z}') \right]_a^b.$$

The boundary conditions imply that this boundary term vanishes at both  $a$  and  $b$ . Hence the operator is self-adjoint. □

### 5 Reality of Eigenvalues

**Theorem 3.** *All eigenvalues of a regular Sturm-Liouville boundary value problem are real.*

*Proof.* Let  $(\lambda, y)$  be an eigenpair. Since  $L$  is self-adjoint,

$$\langle L[y], y \rangle = \langle y, L[y] \rangle.$$

Using the eigenvalue equation,

$$\lambda \int_a^b s(x)|y|^2 dx = \bar{\lambda} \int_a^b s(x)|y|^2 dx.$$

Since  $s(x) > 0$  and  $y \neq 0$ , we conclude  $\lambda = \bar{\lambda}$ . □

## 6 Abel's Formula

**Theorem 4** (Abel's Formula). *Let  $\phi_1$  and  $\phi_2$  be two solutions of (1) corresponding to the same  $\lambda$ . Then*

$$p(x)W(\phi_1, \phi_2)(x) = \text{constant},$$

where

$$W(\phi_1, \phi_2)(x) := \phi_1(x)\phi_2'(x) - \phi_2(x)\phi_1'(x).$$

*Proof.* Multiply the equation for  $\phi_1$  by  $\phi_2$  and that for  $\phi_2$  by  $\phi_1$ . Subtracting yields

$$\frac{d}{dx}(p(x)W(\phi_1, \phi_2)(x)) = 0.$$

Hence the expression is constant. □

## 7 Uniqueness of Eigenfunctions

**Theorem 5.** *For a given eigenvalue  $\lambda$ , any two eigenfunctions are linearly dependent.*

*Proof.* Let  $\phi_1$  and  $\phi_2$  be eigenfunctions corresponding to  $\lambda$ . Both satisfy the boundary conditions, which imply

$$W(\phi_1, \phi_2)(a) = 0.$$

By Abel's formula,  $p(x)W(\phi_1, \phi_2)(x)$  is constant, hence identically zero. Thus  $W(\phi_1, \phi_2) \equiv 0$ , which implies linear dependence:

$$\phi_1 = C\phi_2.$$

□

## Summary

- Distinct eigenfunctions are orthogonal with respect to  $s(x)$ .
- The Sturm–Liouville operator is self-adjoint.
- All eigenvalues are real.
- For a fixed eigenvalue, eigenfunctions are unique up to a multiplicative constant.