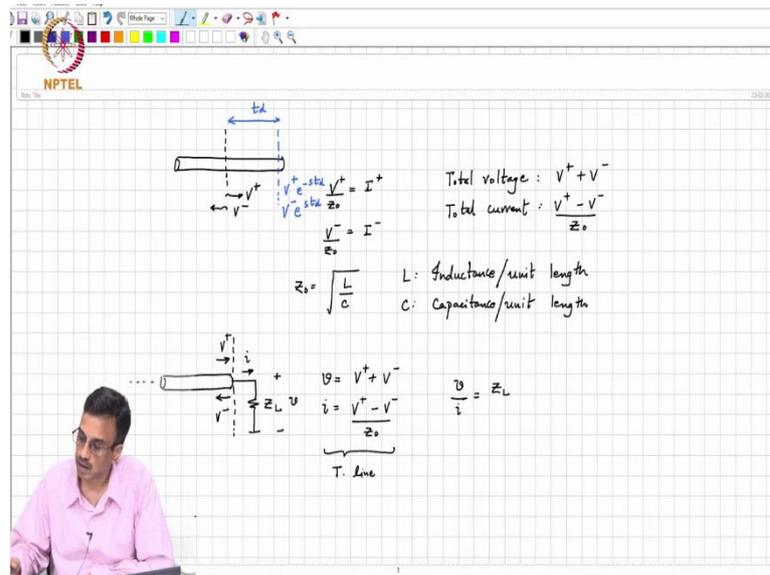


Circuit Analysis for Analog Designers
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Lecture - 45

Transmission line circuit analysis, the reflection coefficient, open- and short-circuited lines

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In the last class, we were talking about transmission lines and the telegraphers equations and we concluded that at any point on ideal transmission line, the voltage can be expressed as the sum of a forward going wave V plus and a backward going wave V minus right. The forward going wave is associated with a current which is I plus and the backward going wave is by the same token I minus.

$$\frac{V^+}{Z_0} = I^+$$

$$\frac{V^-}{Z_0} = I^-$$

The total voltage at any point therefore, is V plus plus V minus.

$$\text{Total voltage: } V^+ + V^-$$

While the total current at any point a line is V plus minus V minus by Z .

$$\text{Total current: } \frac{V^+ - V^-}{Z_0}$$

Recall that Z_0 is the characteristic impedance of the transmission line and is given by a square root of L by C where, L is the inductance per unit length and C is the capacitance of a unit length.

$$Z_0 = \sqrt{\frac{L}{C}}$$

So, as the wave propagates, the positive going wave in other words the wave going towards the right keeps getting delayed. So, for instance if the delay corresponding to this section of the line is t_d , then the positive going wave at the right hand is going to be given by V^+ plus e^{-st_d} and the frequency domain because the wave which is travelling towards you are right is going to get delayed as it travels.

So, and e^{-st_d} is simply the transfer function of an ideal delay with a value t_d . Similarly V^- on the other hand must be advanced when you look at it at the blue boundary. So, that is getting delayed as it moves towards the left and therefore, at the black boundary, it appears V^- alright.

Now, let us see some applications of it turns out that these are very useful relations to know while analyzing circuits with transmission lines. So, in this course we are only going to look at idle transmission lines.

And so for example, consider what happens when you are at the boundary. So, here we do not know what is happening towards the left side the trans the ideal transmission line is terminated by an impedance Z_L and just at the interface we have let us call the positive going wave as V^+ and the negative going wave as V^- and the total voltage is simply $V^+ + V^-$, the total current is $V^+ - V^-$ by Z_0 .

$$v = V^+ + V^-$$

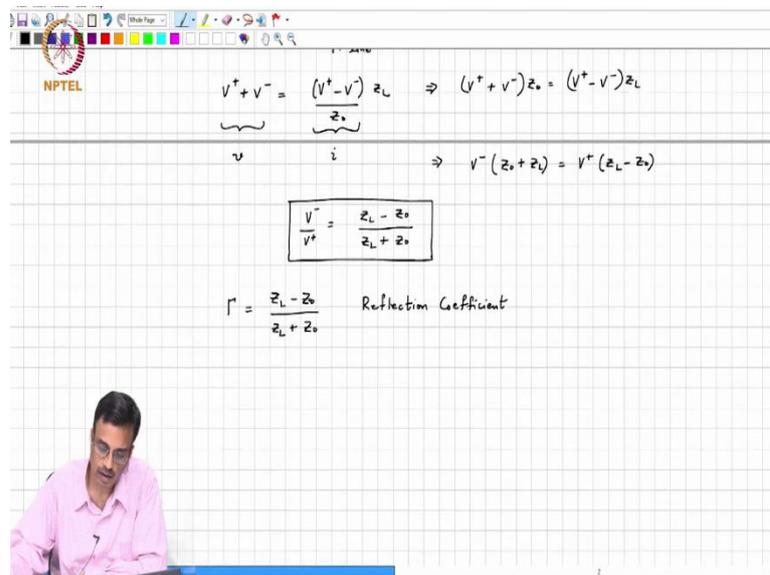
$$i = \frac{V^+ - V^-}{Z_0}$$

So, that is the total current and that is voltage and therefore, this is the constraint imposed by the transmission line and the terminating impedance; however, imposes a constraint due to Ohms law or v by i , its nothing, but Z_L , alright.

$$\frac{v}{i} = Z_L$$

So, this means that we should now be able to figure out what the reverse going wave is in terms of the forward going wave and that is what we will do going forward.

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So, $V^+ + V^-$ must be $V^+ - V^-$ by Z_0 times Z_L .

$$V^+ + V^- = \frac{(V^+ - V^-)}{Z_0} Z_L$$

So, this is the voltage, this on the other hand is the current and therefore, $V^+ + V^-$ into Z_0 is $V^+ - V^-$ into Z_L which means let V^- into $Z_0 + Z_L$ is V^+ into $Z_L - Z_0$. So, we therefore, have V^- over V^+ is nothing, but that $Z_L - Z_0$ by $Z_L + Z_0$.

$$\Rightarrow (V^+ + V^-)Z_0 = (V^+ - V^-)Z_L$$

$$\Rightarrow V^-(Z_0 + Z_L) = V^+(Z_L - Z_0)$$

$$\frac{V^-}{V^+} = \frac{Z_L - Z_0}{Z_L + Z_0}$$

In other words, if you have a transmission line which is terminated at one end by an impedance Z_L , then Ohm's law has to be satisfied as far as the impedance is concerned as far as the terminating impedance Z_L is concerned. However, the forward and backward going waves on the transmission line must. So, what this basically means is that the impedance Z_L imposes a constraint due to Ohm's law; however, the transmission line also imposes constraints in the sense that the quantities on the transmission line must obey the telegrapher's equation.

So, at the interface both these relationships must be satisfied and that is why we have a constraint that the reflected wave is dependent on the terminating impedance and the characteristic impedances of the transmission line of course, its proportional to the amplitude of the incident wave.

So, this quantity Z_L minus Z_0 by Z_L plus Z_0 is called the reflection coefficient and is often denoted by the Greek symbol capital gamma. So, this is the reflection coefficient, alright.

$$\Gamma = \frac{Z_L - Z_0}{Z_L + Z_0}$$

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The slide content is as follows:

$$\frac{V^-}{V^+} = \frac{Z_L - Z_0}{Z_L + Z_0}$$

$$\Gamma = \frac{Z_L - Z_0}{Z_L + Z_0} \quad \text{Reflection Coefficient}$$

If Z_L is freq. dependent: $Z_L(j\omega) \Rightarrow \Gamma(j\omega) = \frac{Z_L(j\omega) - Z_0}{Z_L(j\omega) + Z_0}$

Sanity Check: $Z_L = \text{open circuit}$

$$\Gamma = \frac{Z_L - Z_0}{Z_L + Z_0} = 1$$

Diagram: An open circuit is shown with a voltage source V^+ and a load Z_0 . The current $i=0$ is indicated. The reflection coefficient is calculated as $\frac{V^+ - V^-}{Z_0} = 0 \Rightarrow V^- = V^+ \Rightarrow \Gamma = 1$.

And if Z_L is frequency dependent, in other words Z_L of $j\omega$, then the reflection coefficient is also a function of frequency and is given by Z_L of $j\omega$ minus Z_0 by Z_L of $j\omega$ plus Z_0 , alright.

$$Z_L(j\omega) \Rightarrow \Gamma(j\omega) = \frac{Z_L(j\omega) - Z_0}{Z_L(j\omega) + Z_0}$$

And let us do some sanity checks. The first case we consider is when Z_L is an open circuit. So, in this case Γ , our formula is telling us that Γ which is Z_L minus Z_0 by Z_L plus Z_0 is simply plus 1 because Z_L tends to infinite right.

$$\Gamma = \frac{Z_L - Z_0}{Z_L + Z_0} = 1$$

And why does that make intuitive sense? At the end of the transmission line where we had Z_L being infinity. So, we have an open circuit and the property of an open circuit as we all know is that this current i is 0, if the current i is 0, the current is nothing, but V^+ plus V^- minus Z_0 which must be 0 which can only happen when V^- is equal to V^+ which is equivalent to saying Γ is 1, ok.

$$\frac{V^+ - V^-}{Z_0} = 0$$

$$\Rightarrow V^- = V^+ \Rightarrow \Gamma = 1$$

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The whiteboard content includes:

- Top section:** $Z_L = \text{short circuit}$, $\Gamma = \frac{0 - Z_0}{0 + Z_0} = -1$. A diagram shows a short circuit with $V^+ + V^- = 0$ and $V^- = -V^+ \Rightarrow \Gamma = -1$.
- Middle section:** $Z_L = Z_0$, $\Gamma = \frac{Z_0 - Z_0}{Z_0 + Z_0} = 0$. A diagram shows a load Z_0 with $V^- = 0$.
- Bottom section:** A transmission line diagram with incident wave $V^+ e^{stx}$ and reflected wave $\Gamma V^+ e^{-stx}$. The input impedance is $Z_{in}(s) = ?$. The derivation shows $V = V^+ e^{stx} + \Gamma V^+ e^{-stx}$ and $i = \frac{V^+ e^{stx}}{Z_0} - \frac{\Gamma V^+ e^{-stx}}{Z_0}$. The final result is $Z_{in}(s) = \frac{V}{i} = \frac{V^+ (e^{stx} + \Gamma e^{-stx})}{V^+ (e^{stx} - \Gamma e^{-stx})}$.

Let us now look at another example, Z_L is a short circuit. So, Γ then it is nothing, but $0 - Z_0$ by $0 + Z_0$ which is equal to minus 1

$$\Gamma = \frac{0 - Z_0}{0 + Z_0} = -1$$

And how do we figure this out? This is when the voltage here is 0 and the voltage at the end of the transmission line is V^+ plus V^- which is equal to 0 which means that V^- is equal to minus V^+ , alright. So, so this means that is equivalent to saying that Γ is minus 1.

$$V^+ + V^- = 0$$

$$V^- = -V^+$$

$$\Gamma = -1$$

Now, let us take third case where the impedance is exactly the same as the characteristic impedance of the line

$$Z_L = Z_0$$

And the formula is telling us that Γ is Z_L which is equal to Z_0 minus Z_0 by Z_0 plus Z_0 which is equal to 0 right.

$$\Gamma = \frac{Z_0 - Z_0}{Z_0 + Z_0} = 0$$

And why does this make intuitive sense? This is saying that if you terminate a transmission line from ideal transmission line with its characteristic impedance, the reflected wave is 0 amplitude to the reflected wave is 0.

And why does that make sense? Well, you can think of this as follows, I assume that you had an infinitely long line with the same characteristic impedance. So, this is our line which Z_0 and this goes to infinity. So, this is also Z_0 right. As the wave is concerned this looks like a continuous transmission line which goes to infinity. So, nothing can get reflected back from infinity because it takes an infinitely long time to come back.

So, as far as the transmission line is the backward going wave is concerned that is got to be 0 which of course, is also borne out by the by the reform right. So, intuitively of course, you know you can think of it as the transmission line which is continuing on up to 2 infinity.

Now, if you have a transmission line with the time delay t_d and the characteristic impedance Z_0 and it is terminated by an arbitrary impedance Z_L . The question is what happens, what is the looking in impedance here? Alright. And to do that you know this we just use the techniques that we have developed just now. So, let us assume that this is V^+ , this as we just discussed is ΓV^+ where Γ is $Z_L - Z_0$ by $Z_L + Z_0$, ok write that down here ok.

$$\Gamma = \frac{Z_L - Z_0}{Z_L + Z_0}$$

Now, if the wave at the right end of the transmission line was V^+ , then it must follow that the forward going wave at the left end of the line must be, must have been there a time t_d earlier. So, this must be $V^+ e^{s t_d}$ because after delaying it by t_d it is V^+ at the right end right. In a similar fashion the left going wave at the left end is basically or the backward going way at the left end its nothing, but $\Gamma V^+ e^{-s t_d}$ because the wave going towards the left has gotten delayed by a time t_d .

So, what comment can we make about the voltage at the total voltage at the left end of the line? The voltage is nothing, but the sum of the forward going wave and the backward going wave. So, that is nothing, but $V^+ e^{s t_d} + \Gamma V^+ e^{-s t_d}$, the current is nothing, but $V^+ e^{s t_d} / Z_0 - \Gamma V^+ e^{-s t_d} / Z_0$.

$$v = V^+ e^{s t_d} + \Gamma V^+ e^{-s t_d}$$

$$i = \frac{V^+ e^{s t_d}}{Z_0} - \frac{\Gamma V^+ e^{-s t_d}}{Z_0}$$

So, the impedance Z_{in} is simply the ratio of the voltage to the current at the left end of the line and that is nothing, but $V^+ e^{s t_d} + \Gamma V^+ e^{-s t_d}$ divided by $V^+ e^{s t_d} / Z_0 - \Gamma V^+ e^{-s t_d} / Z_0$.

$$Z_{in}(s) = \frac{v}{i} = \frac{V^+(e^{st_d} + \Gamma e^{-st_d})}{V^+(e^{st_d} - \Gamma e^{-st_d})} Z_0$$

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The slide shows the following derivations:

$$Z_{in}(s) = ?$$

$$v = V^+ e^{st_d} + \Gamma V^+ e^{-st_d}$$

$$i = \frac{V^+ e^{st_d}}{Z_0} - \frac{\Gamma V^+ e^{-st_d}}{Z_0}$$

$$Z_{in}(s) = \frac{v}{i} = \frac{V^+ (e^{st_d} + \Gamma e^{-st_d})}{V^+ (e^{st_d} - \Gamma e^{-st_d})} Z_0$$

$$Z_{in}(s) = Z_0 \left(\frac{e^{st_d} + \Gamma e^{-st_d}}{e^{st_d} - \Gamma e^{-st_d}} \right)$$

$$Z_{in}(j\omega) = Z_0 \left(\frac{e^{j\omega t_d} + \Gamma e^{-j\omega t_d}}{e^{j\omega t_d} - \Gamma e^{-j\omega t_d}} \right)$$

Special case: $\omega t_d = \pi/2 \Rightarrow 2\pi f t_d = \frac{\pi}{2} \Rightarrow t_d = \frac{1}{4f}$
quarter wave line

So, the impedance is simply given by sorry e this times Z, impedance is simply given by there in a s is Z naught times e to the s t d plus gamma e to the minus s t d by e to the s t d minus gamma minus s t d ok.

$$Z_{in}(s) = Z_0 \left(\frac{e^{st_d} + \Gamma e^{-st_d}}{e^{st_d} - \Gamma e^{-st_d}} \right)$$

And all that we have the, all that we have used to get this relationship is simply observe two things. One at the boundary condition Ohm's law must be satisfied along with the telegraphers equations and at the along the line the forward wave propagates towards the right, the backward wave propagates towards the left and they have the same delay.

At the left end of the line the total voltage is the sum of the forward and the backward going waves and similarly the current is the sum of the current due to the forward going wave which is in the right and the current that is due to the backward going wave which is in the opposite direction. So, the current has to be the two currents have to be subtracted from each other.

So, and the impedance is simply the ratio of the voltage to be cut and with that we get this expression. Now, let us take a look at this expression you know not only to get familiar

with it, but also a useful practical case. A special case occurs remember that when you are doing this the sinusoidal steady state then this now all these become s equal to j omega.

So, this is nothing, but Z_{in} is $Z_0 \frac{e^{j\omega t_d} + \Gamma e^{-j\omega t_d}}{e^{j\omega t_d} - \Gamma e^{-j\omega t_d}}$

$$Z_{in}(j\omega) = Z_0 \left(\frac{e^{j\omega t_d} + \Gamma e^{-j\omega t_d}}{e^{j\omega t_d} - \Gamma e^{-j\omega t_d}} \right)$$

A special case occurs when ωt_d is $\pi/2$, alright.

$$\omega t_d = \frac{\pi}{2}$$

So, what does this, what does it actually mean? Remember right ω is $2\pi f$ times t_d is $\pi/2$ which basically means that t_d is one-fourth of $1/f$

$$\Rightarrow 2\pi f t_d = \frac{\pi}{2} \Rightarrow t_d = \frac{1}{4f}$$

And let us stare at this and try and see if we can interpret this physically $1/f$ is nothing but the time period of the sinusoidal right the input sinusoidal is at a frequency f , $1/f$ therefore, is a time period and in one time period of the wave a wave will travel one wavelength. So, this t_d if it is chosen such that ωt_d is 90 degrees $\pi/2$ radians intuitively or physically that corresponds to the delay of the transmission line being equal to quarter wave line. So, this is what is called a quarter wave line, ok.

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Special case: $\omega t_d = \pi/2 \Rightarrow 2\pi f t_d = \frac{\pi}{2} \Rightarrow t_d = \frac{1}{4f}$
 quarter wave line
 $e^{j\omega t_d} = j$

$$Z_{in}(j\omega) = Z_0 \left(\frac{j - \Gamma j}{j + \Gamma j} \right) = Z_0 \frac{(1 - \Gamma)}{(1 + \Gamma)}$$

$$\Gamma = \frac{Z_L - Z_0}{Z_L + Z_0} \Rightarrow Z_0 \frac{1 - \frac{Z_L - Z_0}{Z_L + Z_0}}{1 + \frac{Z_L - Z_0}{Z_L + Z_0}} = Z_0 \frac{2Z_0}{2Z_L}$$

$$\Rightarrow Z_{in}(j\omega) = \frac{Z_0^2}{Z_L} \quad \text{Impedance inversion}$$

$Z_L = \text{open} \Rightarrow Z_{in} = 0$	} at a particular frequency
$= \text{short} \Rightarrow Z_{in} = \infty$	
$= jX \Rightarrow -j \left(\frac{Z_0^2}{X} \right)$	
inductive Capacitive	

Now, so Z_{in} or $j\omega$ if this is the case then e to the j by e to the $j\omega t_d$ is simply given by $\cos \frac{\pi}{2} + j \sin \frac{\pi}{2}$ which is simply nothing but j , alright. So, Z_{in} of $j\omega$ is nothing but j . So, I mean either $j\omega t_d$ is nothing, but j . So, Z_{in} is nothing but Z_0 times j plus Γ times j . So, that becomes and this becomes j plus Γ which is nothing, but Z_0 times $1 - \Gamma$ by $1 + \Gamma$ and recall that Γ is nothing but $Z_L - Z_0$ by $Z_L + Z_0$.

So, which basically means that looking in impedance is nothing but Z_0 times that Z_L plus Z_0 $1 - Z_L - Z_0$ by $Z_L + Z_0$ divided by $1 + Z_L - Z_0$ by $Z_L + Z_0$ ok. And this is basically Z_0 times $2Z_0$ divided by $2Z_L$ therefore, Z_{in} of $j\omega$ is Z_0^2 by Z_L , alright.

$$e^{j\omega t_d} = j$$

$$Z_{in}(j\omega) = Z_0 \left(\frac{j - \Gamma j}{j + \Gamma j} \right) = Z_0 \frac{(1 - \Gamma)}{(1 + \Gamma)}$$

$$\Gamma = \frac{Z_L - Z_0}{Z_L + Z_0} \Rightarrow Z_0 \frac{1 - \frac{Z_L - Z_0}{Z_L + Z_0}}{1 + \frac{Z_L - Z_0}{Z_L + Z_0}} = Z_0 \frac{2Z_0}{2Z_L}$$

$$Z_{in}(j\omega) = \frac{Z_0^2}{Z_L}$$

So, as you can see the impedance looking on the left is an inverted version of the terminating impedance. So, and, so if you put a capacitor it will look like an inductor, an open circuit. So, if Z_L is open, Z_{in} will look like a short, if that Z_L is a short Z_{in} will look like an open circuit, if Z_L is inductive, then it will look like, so, this is inductive it will look like minus j times Z_0^2 over X , $\left(-j \left(\frac{Z_0^2}{X}\right)\right)$ and therefore, this those capacitive.

With a key point to note is that all this only happens at that frequency ω which makes the length of the line equal to quarter wave.