

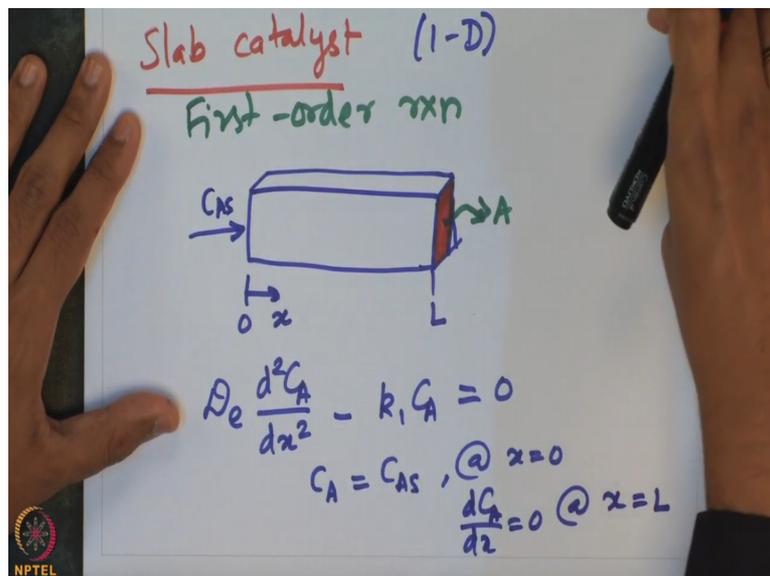
Chemical Reaction Engineering - II
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Module - 5
Lecture - 22
Internal Effectiveness Factor II

In the last lecture we derived the relationship between the internal effectiveness factor and the Thiele modulus for a first order reaction occurring in a spherical catalyst pellet. We also observed that such a relationship can actually be obtained for any order of a reaction that is happening in the catalyst pellet. Now, we also noted that the catalyst pellet can be of different shapes.

For instance, it can be a cylindrical pellet or it can have a Cartesian coordinate, where it is a slab geometry. So, in today's lecture, we will look at how to find out the relationship between the effectiveness factor and Thiele modulus, for a slab geometry, for a first order reaction.

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So, suppose we have a slab. And let us say that I define my positive direction x . And let us say that the slab goes from 0 to L . Length of the slab is L . And the species is actually diffusing from this direction. And suppose, let us say I seal the other end of the catalyst, so that nothing is actually leaving this surface. And if I assume that in all the other dimensions, the concentration is actually uniform.

Let us look at a one-dimensional case. So, let us look at a 1 D slab. And if I assume that the cross-sectional area through which the species is diffusing is actually = A. So, that is the cross-sectional area. So, the mass balance for this system is essentially given by; one can actually write the shell balances. And you can find that the balance is essentially given by diffusivity into d square C A by d x square, - k 1 into C A is = 0.

What are the boundary conditions? So, concentration of the species will be = the surface concentration @ x = 0. Essentially, the species at this location is at the surface concentration. And because nothing is leaving this, from this end at L, so the next boundary conditional will be d C A by d x, essentially is = 0, @ x = L. So, we can solve this model in exactly the same way we did for a spherical catalyst by first non-dimensionalizing the model equation. So, let us do that exercise. So now, we introduce the dimensionless quantities.

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$$\psi = \frac{C_A}{C_{As}}, \quad \lambda = \frac{x}{L}$$

$$\psi = \frac{C_A}{C_{As}} = 1 \quad @ \quad \lambda = 0$$

$$\frac{d\psi}{d\lambda} = 0 \quad @ \quad \lambda = 1$$

$$D_e \frac{C_{As}}{L^2} \frac{d^2\psi}{d\lambda^2} - k_1 C_{As} \psi = 0$$

$$\frac{d^2\psi}{d\lambda^2} - \phi_1^2 \psi = 0$$

Is = C A by C A S. And lambda is = x by L. So clearly, lambda goes between 0 and 1. And so now, the boundary conditions will be, psi = C A by C A S, is = 1 @ lambda = 0. And d psi by d lambda is essentially = 0 @ lambda = 1. Now, if we now substitute the dimensionless quantities in the model equation, we will see that diffusivity D e into C A S into d square psi, divided by R square into d lambda square, - k 1 into, C A S into psi is = 0.

Now, we can cancel out the like terms. And then, we can pull in the D e, we can basically take this D e by R square into the second term. And so, we will essentially, the model equation, dimensionless model equation will be d square psi by d lambda square, - phi 1 square, into psi = 0. So, that is the model equation, which will be subject to these 2 boundary

conditions, where $\psi = 1$ and $\lambda = 0$. And $d\psi/d\lambda = 0$, at $\lambda = 1$. So, let us solve this model first. So, what is the general solution?

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The image shows a whiteboard with the following handwritten mathematical work:

$$\psi = A e^{\phi \lambda} + B e^{-\phi \lambda}$$

$$\phi = \sqrt{\frac{k_1}{D_e}} L$$

$$\psi = 1 \text{ @ } \lambda = 0$$

$$A + B = 1 \Rightarrow A = 1 - B$$

$$\frac{d\psi}{d\lambda} = 0 \text{ @ } \lambda = 1$$

$$\left. \frac{d\psi}{d\lambda} \right|_{\lambda=1} = [A\phi_1 e^{\lambda\phi_1} - B\phi_1 e^{-\phi_1\lambda}]_{\lambda=1}$$

$$= A\phi_1 e^{\phi_1} - B\phi_1 e^{-\phi_1} = 0$$

Additional notes on the right side of the board:

$$A = 1 - B$$

$$A e^{\phi_1} - B e^{-\phi_1} = 0$$

General solution is, $\psi = A e^{\phi x} + B e^{-\phi \lambda}$. So now, ϕ is essentially = square root of k_1 by D_e into L . Note that this is not radius. This is L which is the length of the slab. It is L square. So, $\phi = \sqrt{k_1 / D_e} L$. And so now, we can substitute the boundary condition. The first boundary condition is $\psi = 1$ @ $\lambda = 0$.

So, substituting that, we will see that $A + B = 1$; which means that $A = 1 - B$. So, when I substitute $\lambda = 0$, $e^0 = 1$. And this is 1. So, therefore $A + B = 1$. And $A = 1 - B$. Now, I substitute the second boundary condition. $d\psi/d\lambda = 0$ @ $\lambda = 1$. And so, $d\psi/d\lambda$ is essentially =; that is = $A\phi_1 e^{\phi_1} - B\phi_1 e^{-\phi_1} = 0$.

So, if I now put subscript 1 to represent the first order reaction. So, $A\phi_1 e^{\phi_1} - B\phi_1 e^{-\phi_1} = 0$, evaluated at $\lambda = 1$. So now, if I put $\lambda = 1$; this is essentially becomes $A\phi_1 e^{\phi_1} - B\phi_1 e^{-\phi_1} = 0$. And that is essentially = 0. So, from here, we get 2 conditions on A and B .

One is that $A = 1 - B$. And the other one is, $A\phi_1 e^{\phi_1} - B\phi_1 e^{-\phi_1} = 0$. Because this is thiele modulus. And thiele modulus clearly cannot be 0. So therefore, it is

the other product which is actually = 0. So, from this we can actually find out the constants A and B. Let us do that. So, the solution of these 2 equations will essentially be $A = e^{-\phi_1} / (e^{\phi_1} + e^{-\phi_1})$.

And B is equal $e^{\phi_1} / (e^{\phi_1} + e^{-\phi_1})$. So, substituting this back into the solution, you will find that $\psi = e^{-\phi_1} \cosh(\lambda z) + e^{\phi_1} \sinh(\lambda z) / (e^{\phi_1} + e^{-\phi_1})$. And so, rewriting this by taking this into, by joining these 2 exponentials, we will see essentially that this is $\cosh(\phi_1(1-\lambda)) / \cosh(\phi_1)$.

And this is, if I now multiply and divide by a factor 2. This is nothing but, this is $\cosh(\phi_1(1-\lambda)) / \cosh(\phi_1)$. So, the solution ψ which is the dimensionless concentration as a function of position, is essentially given by this ratio, $\cosh(\phi_1(1-\lambda)) / \cosh(\phi_1)$. So, once again we can find out the internal effectiveness factor. How do we do that? We do the same exercise as we did for spherical catalyst;

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The image shows a whiteboard with handwritten mathematical derivations. The first equation is $\eta = \frac{-r_A}{-r_{A,s}} = \frac{-D_e \frac{dC_A}{dz} \big|_{z=0} \cdot A}{k_1 C_{A,s} A \cdot L}$. The second equation is $-D_e \frac{dC_A}{dz} \big|_{z=0} = -D_e C_{A,s} \frac{d\psi}{dz} \big|_{z=0}$. The third equation is $\psi = \frac{\cosh(\phi_1(1-z))}{\cosh(\phi_1)}$. The fourth equation is $\frac{d\psi}{dz} = \frac{-\phi_1 \sinh(\phi_1(1-z))}{\cosh(\phi_1)}$.

Where we say that η is essentially $-r_A$ which is the actual rate at which the reaction is happening, divided by the rate as though the conditions everywhere inside the slab is that of the surface conditions or the concentration at $x = 0$. That is, $C_{A,s}$ which is the concentration at $x = 0$. And this is nothing but $-D_e \frac{dC_A}{dz} \big|_{z=0}$ multiplied by the cross-sectional area through which the diffusion is happening, divided by $k_1 C_{A,s}$, which is

the intrinsic kinetics, multiplied by the volume of the slab which is area into the length, cross section area into length.

So, we need to estimate this dC_A by dx at $x = 0$. How do we do that? We know the solution. So, from the solution, we can actually find out what this one is. So, $-D_e dC_A$ by dx at $x = 0$. By introducing the dimensionless quantities, we will see that this is $= -D_e C_{A,S}$ into $d\psi$, divided by L into $d\lambda$. And that should be, $d\psi$ by $d\lambda$. That should be evaluated at $\lambda = 0$. So, this is $d\psi$. $d\psi$ by $d\lambda$. And that should be evaluated at $\lambda = 0$.

But we know what is the relationship between ψ and λ . So, ψ is essentially $= \cos$ hyperbolic of ϕ_1 , $1 - \lambda$, divided by \cos hyperbolic of ϕ_1 . So, from this we can find out what is the first derivative. So, the first derivative $d\psi$ by $d\lambda$ is essentially given by $-\phi_1 \cos$ hyperbolic $1 - \lambda$, $-\phi_1 \sin$ hyperbolic ϕ_1 into $1 - \lambda$, divided by \cos hyperbolic of ϕ_1 . Now, we need to evaluate this at $\lambda = 0$. So, let us do that.

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The whiteboard shows the following steps:

$$\left. \frac{d\psi}{d\lambda} \right|_{\lambda=0} = -\phi_1 \frac{\sinh \phi_1}{\cosh \phi_1} = -\phi_1 \tanh \phi_1$$

$$-D_e \left. \frac{dC_A}{dx} \right|_{x=0} = -D_e \frac{C_{A,S}}{L} \left. \frac{d\psi}{d\lambda} \right|_{\lambda=0}$$

$$= D_e \frac{C_{A,S}}{L} \phi_1 \tanh \phi_1$$

$$\eta = \frac{D_e C_{A,S} \phi_1 \tanh \phi_1 \cdot A}{k_1 C_{A,S} \cdot A \cdot L}$$

So, $d\psi$ by $d\lambda$, at $\lambda = 0$ is essentially given by $-\phi_1$ into \sin hyperbolic of ϕ_1 divided by \cos hyperbolic of ϕ_1 . Now, this is $= -\phi_1 \tan$ hyperbolic of ϕ_1 . So, this we can substitute back into the effectiveness factor equation which we just solved. We can substitute it back here. And so, that we will essentially be, so $-D_e$ into dC_A by dx at $x = 0$ is $=, -D_e C_{A,S}$ by L into $d\psi$ by $d\lambda$ at $\lambda = 0$.

And we have found out what is $d\psi$ by $d\lambda$. So, this is essentially $D_e C_{A,S}$ by L into $\phi_1 \tan$ hyperbolic of ϕ_1 . And this we can substitute back and we can find out that, η

which is the effectiveness factor is essentially given by $D_e C A S$ divided, into $\phi_1 \tanh \phi_1$ hyperbolic of ϕ_1 , divided by L multiplied by the area. The whole divided by $k_1 C A S$ into area into L .

So, we can now cancel out the like terms. We can cancel the cross-sectional area, the surface concentration. And so, what we are left with is essentially $D_e \phi_1$ and k_1 and L square. So, let us write this.

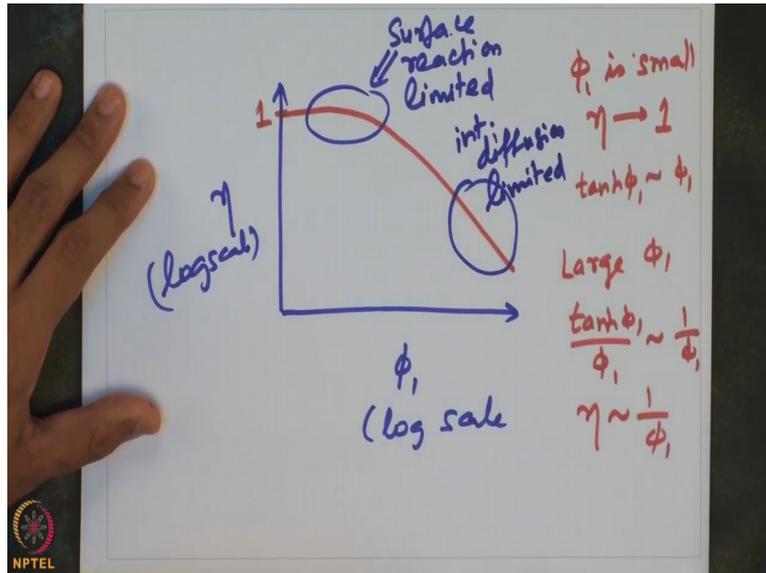
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The image shows a hand-drawn derivation on a whiteboard. The first equation is $\eta = \frac{\phi_1 \tanh \phi_1}{(k_1 L^2 / D_e)} \Rightarrow \phi_1^2$. The second equation is $\eta = \frac{\phi_1 \tanh \phi_1}{\phi_1^2} = \frac{\tanh \phi_1}{\phi_1}$. The final result is boxed: $\eta = \frac{\tanh \phi_1}{\phi_1}$. There are also some small logos and text in the top right corner of the whiteboard image, including 'Destined to rule the land...', 'www.pvcinema.com', 'Aquaman - Now Showing at...', 'PVR', and 'Set'.

So, η is essentially given by $\phi_1 \tanh \phi_1$, divided by $k_1 L^2$ by D_e . Now, we know that $k_1 L^2$ by D_e is nothing but ϕ_1^2 . This is nothing but ϕ_1^2 . So, substituting that we find that η is essentially = $\phi_1 \tanh \phi_1$, divided by ϕ_1^2 . Or that is = $\tanh \phi_1$ by ϕ_1 . So, for a slab geometry, we find that η is essentially = $\tanh \phi_1$ divided by ϕ_1 .

That is the relationship between the internal effectiveness factor and the thiele modulus for a slab geometry. Now we can plot this, once again we can plot this η versus ϕ_1 . And we will see that the profile that we get is very similar to what was observed for a spherical catalyst pellet.

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So, if I now plot ϕ_1 versus η , once again log scale. I use log scale here. And clearly, when ϕ_1 is very small, η essentially goes to 1, because \tanh hyperbolic of ϕ_1 essentially goes as ϕ_1 . So, when ϕ_1 is small, so \tanh hyperbolic of ϕ_1 essentially goes as ϕ_1 . And so, η essentially tends to 1. And for large values of ϕ_1 , for large ϕ_1 , \tanh hyperbolic of ϕ_1 by ϕ_1 , essentially scales as 1 by ϕ_1 .

So therefore, η goes as 1 by ϕ_1 for large values of ϕ_1 . So clearly, we can draw this η versus ϕ_1 , η versus ϕ_1 . And recall that, for surface limited case, for small ϕ_1 η goes to 1, similar to what we saw in the spherical catalyst pellet. And also for large ϕ_1 , where it is diffusion limited conditions, η scales as 1 over ϕ_1 , which is the similar for, similar to what we saw in a spherical catalyst pellet.

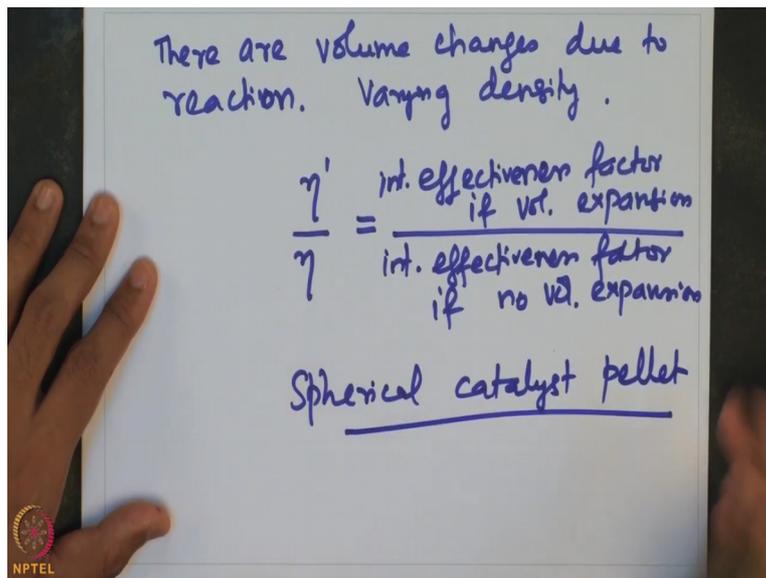
Except that that⁰ we had a 3 by ϕ_1 . So, the scaling factor or the slope of the line, for large values of ϕ_1 is only different. So, now we can draw the sketch the thiele modulus versus effectiveness factor graph. So, this is, starts from 1. And so, this is the, this region is now the surface reaction limited. And this region is now the internal diffusion limited. So, anywhere in between, it could be both surface reaction and internal diffusion that actually contributes to the overall rate at which the reaction can happen.

The same procedure can actually be followed for finding out the η versus ϕ_1 relationship for a cylindrical catalyst pellet. And in fact, you will see that the relationship between η and ϕ_1 for a cylindrical catalyst pellet is also very similar to what we saw for a slab or what we saw for a spherical catalyst pellet. So, what is remarkable is that the, irrespective of the shape

of the catalyst pellet, the by and large, the functional relationship or the dependency of eta versus phi is actually very similar.

And in fact, for large values of phi, irrespective of the geometry that we consider, the effectiveness factor essentially scales as 1 over thiele modulus. And that actually is quite an interesting observation. Now, so far what we considered is the case where the density of the system is constant.

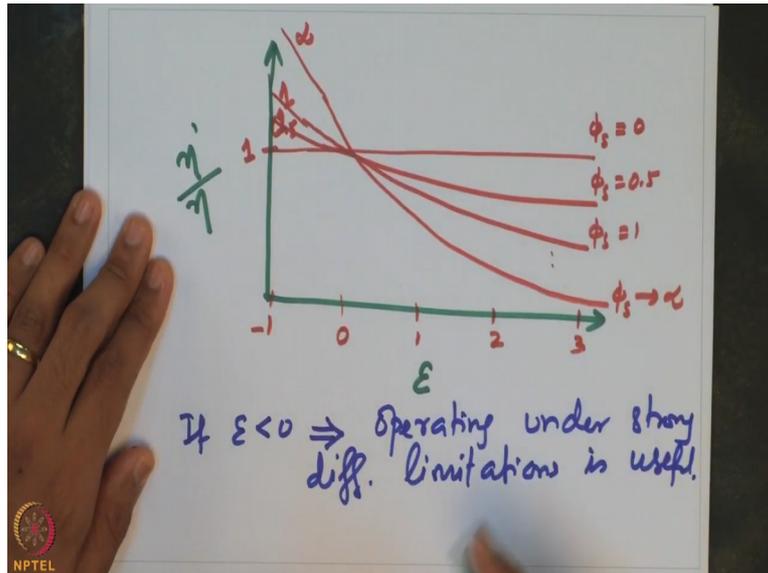
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Suppose, let us say that there are volume changes due to reaction. Let us say that there are volume changes due to reaction. And if, let us say that, which basically means that we are looking at varying density systems. So now, if you want to understand how the eta versus phi actually behaves when the density is changing. So, we can actually try to relate it to what is called as the ratio of the effectiveness factor, if there is volume changes, divided by eta if there is no volume changes.

So, we can try to find the ratio of this, which is basically the effectiveness, internal effectiveness factor, if volume expansion, divided by the corresponding internal effectiveness factor if there is no, if no volume expansion. So, let us consider the case for a spherical catalyst pellet. Now in fact, for a spherical catalyst pellet, this has been worked out and it is essentially captured by this graph here.

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Where we can plot, eta prime by eta, divided by the expansion coefficient epsilon. And let us say that there is no volume change or 0 expansion that corresponds to this location. And let us say - 1 is here, + 1 is here, 2 is here and 3. So, this essentially says what is the extent to which the molar changes happen because of the reaction. And if, let us say that, if there are no volume change effects, then eta prime by eta is essentially same, which means, 1 corresponds to eta prime being = eta.

So, suppose let us say this is the, this is the location where the diffusional, where this essentially says that eta is = eta prime, for all values of this. And suppose, let us say this corresponds to the thiele modulus. Let us say the thiele modulus of approximately 0. That is, no diffusional limitations. Which means that very small or 0 thiele modulus, where every, for every case it is essentially the surface reaction which is controlling.

In that case eta prime is essentially = eta. And now, if I plot this curve for different values of thiele modulus. We will see that, for phi; so, this is for phi s = 0.5. And then, so this is for a different value of phi s. We will say phi s is = 1. And so on and so forth. And typically for infinite, very large phi s, this is the kind of profile. So, this is phi s corresponding to 1. And this is let us say this is 0.5. And let us say this is 1.

And this corresponds to tending to infinity. Now, this graph actually provides a useful insight in terms of design. That, if epsilon is less than 0, then operating under strong diffusional limitation is useful. So, classically for the cases that we have seen before, we saw that the

effectiveness factor is maximum when the reaction is actually conducted at a surface reaction conditions. And not the diffusional limitations.

Because effectiveness factor is approximately = 1 when the surface reaction is controlling and which is the maximum that can be obtained for all these cases. On the other hand, if there are density or volume changes, then we will see that when the, when there is negative volume change, that is the epsilon is -1, then if you operated the strong diffusional limitations, then we are able to get a much higher effectiveness factor.

Eta prime is actually much larger than eta if there was no volume expansion. So clearly, that if there is volume shrinking because of the reaction, then one would actually want to conduct the reaction under strong diffusional limitations. And that would actually help in increasing the overall rate at which the reaction can happen, which actually strongly affects the productivity. So, let us take an example. Let us look at some numbers.

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$$D_{AB} \frac{d^2 C_A}{dx^2} - k = 0$$

$$\psi = \frac{C_A}{C_{A0}}, \quad \lambda = \frac{x}{L}$$

$$\frac{d^2 \psi}{d \lambda^2} - \frac{2 k L^2}{2 D_{AB} C_{A0}} = 0$$

$$\phi^2 = \frac{k L^2}{2 D_{AB} C_{A0}}$$

$$\boxed{\frac{d^2 \psi}{d \lambda^2} - 2 \phi^2 = 0}$$

Suppose we have, so this problem has been taken from Fogler, 4th edition, page number 824. So, suppose there are hydrogels which are present, which contains hydrogel containing cells. And in fact, these days it is being considered as an alternative for tissue replacement; where you have cells of the host species, is actually embedded in the hydrogels and can we see if that can actually replace the cartilages for instance.

So, if these cells have to survive in normal post conditions. They do not have blood vessels. And, but the oxygen has to penetrate for cells to survive. So, one can actually post this as a

diffusion reaction problem and see how diffusional limitations could play a role here. So, suppose I consider that there is a hydrogel which is present here. And I pass oxygen from both sides. And if we assume that there is a midpoint that is present here.

If I called this as, let us say $x = 0$; where is a, there is a symmetric point that is present here. And now, I can write my, and let us say that the length is $x = L$ here and $x = -L$. And I can now write the model equations to find out what should be the length that would actually that would actually support the, you know, diffusion, support the diffusion of oxygen all the way to the end of the, or to midpoint of the hydrogel; in such a way that sufficient quantity of oxygen is always present at that location.

Now, suppose if we assume that the oxygen consumption is a 0 order reaction, we assume that oxygen consumption is actually occurring through a 0 order reaction. We can now write a mass balance between 0 and L, assuming that it is symmetric. So, if we write a balance between 0 and L; it is as good as writing a model between 0 and $-L$. So, let us write the model for this half section. So, if $D A B$ is equimolar counter diffusion of oxygen.

That multiplied by $d^2 C_A$ by $d x^2$, $-k = 0$. Because we assume it is a 0 order reaction. Once again, we can introduce the dimensionless quantity. ψ is $= C_A$ by $C_{A S}$. If we assume that the concentration of oxygen at the location where it enters the gel is actually $C_{A S}$, the concentration here is C_A . Then we can actually rewrite the model in terms of the dimensionless quantities, where λ is $= x$ by L .

And we can rewrite the model as $d^2 \psi$ by $d \lambda^2$, $-k L^2$ by $D A B$ into $C_{A S} = 0$. Now, let us introduce a factor of 2 which actually will help in solving the model equation. So, we will introduce a, multiply and divide the second term by a factor of 2. And so, we can now define ϕ^2 as k into L^2 by 2 times $D A B$ into $C_{A S}$. And so, the model equation now boils down to $d^2 \psi$ by $d \lambda^2$.

That is, $-2 \phi^2 = 0$. Now, what is the solution of this equation? It is very easy to find the solution of this equation.

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$$\psi = \phi^2 \lambda^2 - 2\phi^2 \lambda + K_2$$

$$\boxed{\psi = \phi^2 \lambda (\lambda - 2) + 1}$$

All values of ϕ^2 ?

Suppose O_2 is completely consumed
at $\lambda = 1$ \downarrow
 $\psi = 0$

$$0 = \phi^2 (1)(-1) + 1 \Rightarrow \underline{\underline{\phi^2 \leq 1}}$$

The solution is essentially given by, $\psi = \phi^2 \lambda^2 - 2\phi^2 \lambda + K_2$. And so, from using the boundary conditions we can find out what is this constant K_2 . And so, essentially the solution is given by $\psi = \phi^2 \lambda (\lambda - 2) + 1$. So, this is the relationship between concentration of oxygen at any location inside the gel, as a function of position λ .

Now, is this valid under all situations? Is it valid for all values of ϕ ? Is it valid for all values of ϕ^2 ? So, it turns out that the solution, in fact the model itself is not valid for all values of ϕ^2 . So, because the oxygen is being consumed through a 0 order reaction; somewhere, let us say that the oxygen concentration becomes 0 or oxygen is completely exhausted.

So, let us find out what is this ϕ^2 which will actually permit that what is this location λ and corresponding ϕ^2 , which actually permits complete consumption of the oxygen species. Suppose O_2 is completely consumed at $\lambda = 1$, which is the domain of the hydrogel which we have considered. We can substitute this. So, that is, completely consumed, that is $\psi = 0$.

So, we can substitute that and see at 0 is $\psi = \phi^2 (1)(-1) + 1$, which means that the solution is valid only when ϕ^2 is actually less than or $= 1$. So, the model is not valid at every location. The solution is not valid at, for all values of ϕ^2 . It is valid only when ϕ^2 is actually less than or $= 1$. So, let us summarise what we have done now.

We have actually looked at the thiele modulus and effectiveness factor relationship for a slab geometry. We also looked at, what is the effect of having volume changes, where the effectiveness factor can actually be much larger than 1, depending upon the thiele modulus and if the volume is actually shrinking because of the reaction. We looked at an example problem where we looked at oxygen diffusion through hydrogels containing cells.

And we looked at the solution and found out that the solution is not valid for all values of ϕ square. And what we will see in the next class is, we will put some numbers to these and see what should be the size of the, what should be the thickness of the gel or what should be the value of L that actually is mimicking the realistic situation. Thank you.